



UKAEA

Plasma Scenario Development in the Spherical Tokamak for Energy Production (STEP)

**Ken McClements for the STEP Plasma, Control and Heating & Current Drive Team
21st International Spherical Torus Workshop (ISTW 2022), November 1 2022**

Thanks to the STEP Plasma team & Collaborators ...

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Site for STEP prototype fusion energy plant selected

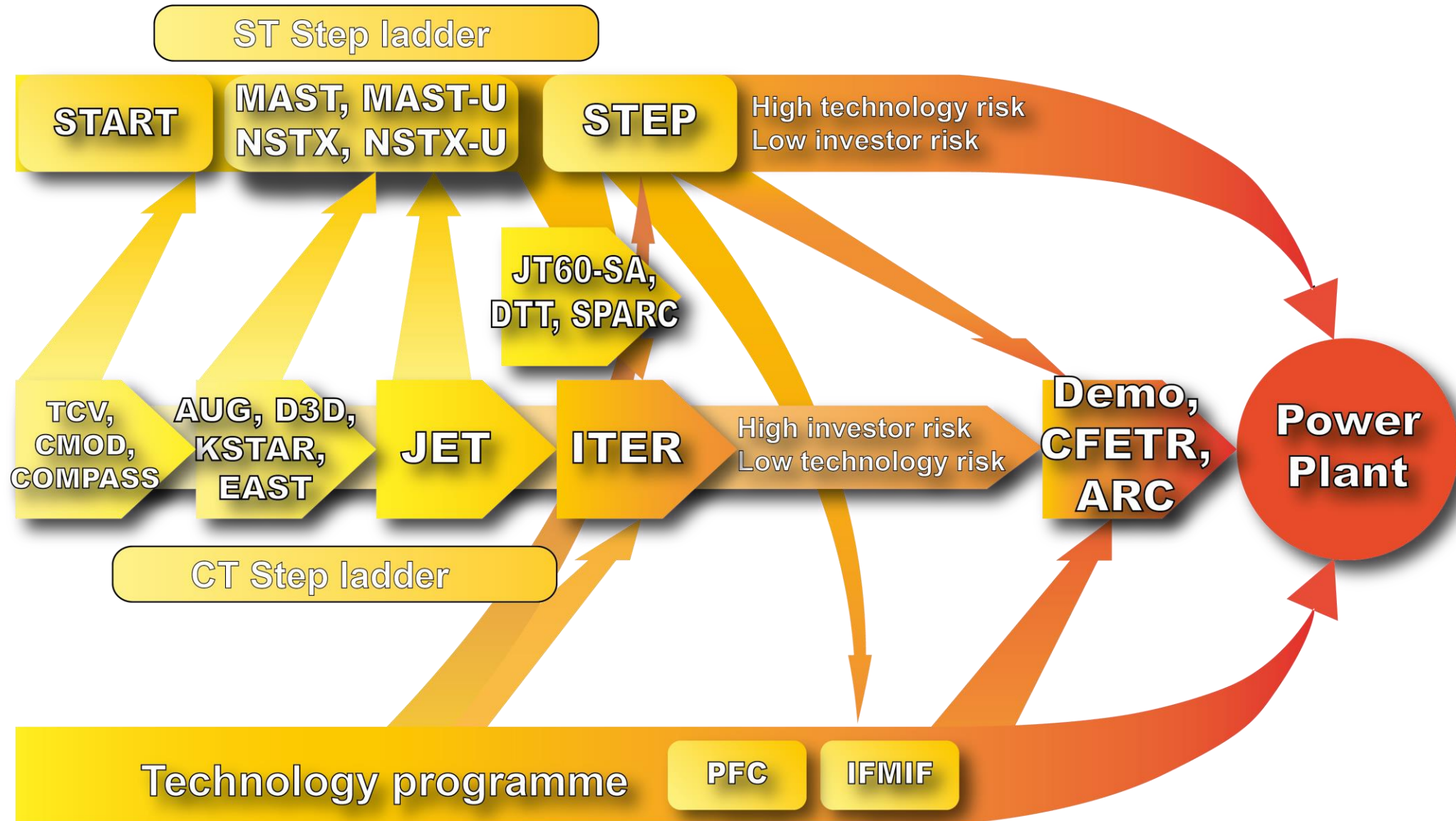
- STEP prototype will be built at West Burton UK, ~200km from Culham Science Centre
- Target completion date: 2040
- Key aims:
 - Fusion power generation on GW scale
 - Net electric power
 - Tritium self-sufficiency



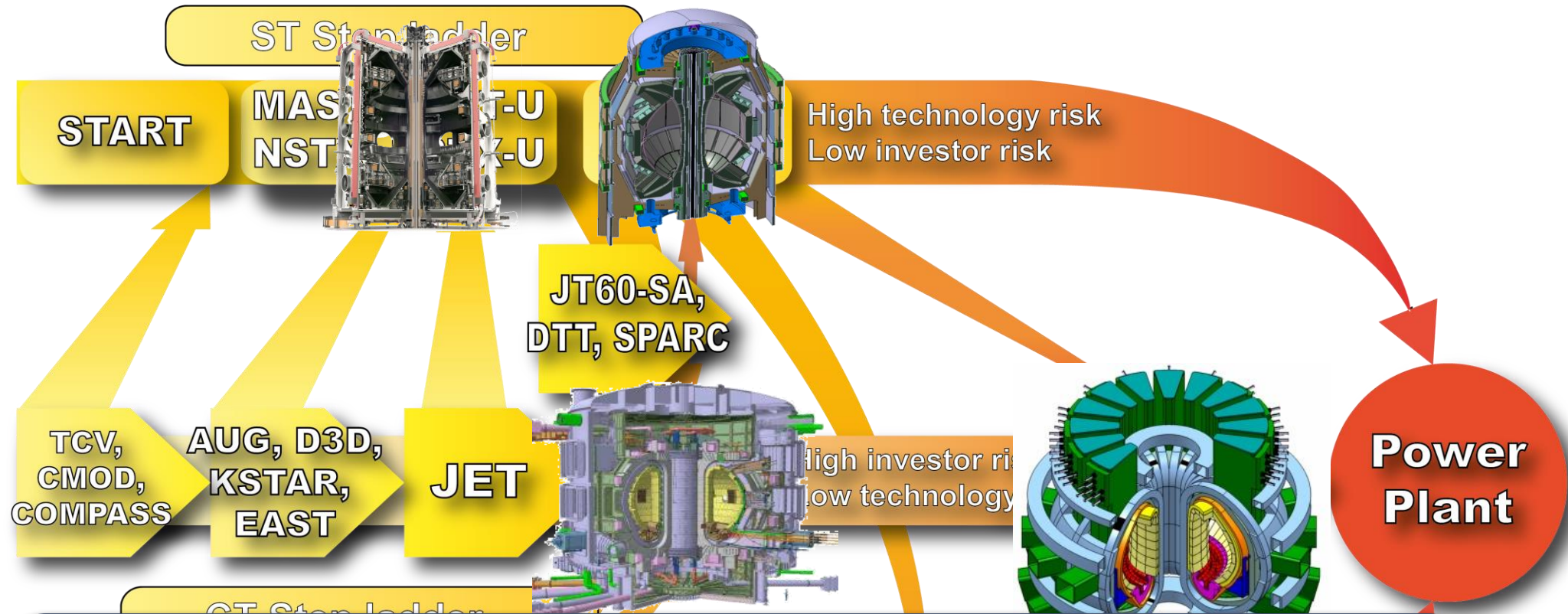
Outline

- Role of STEP in path to power plant
- Plasma modelling workflow
- Flat-top operating points
- Divertor design
- Stability control
- EC & EBW current drive
- Core confinement
- α -particle confinement
- Plasma start-up

Role of STEP in path to power plant



Role of STEP in path to power plant



Deliver UK prototype fusion energy plant targeting 2040 & path to commercial viability of fusion

T1: Concept
2019 - 2024

Tranche 2: Design
2024 - 2032

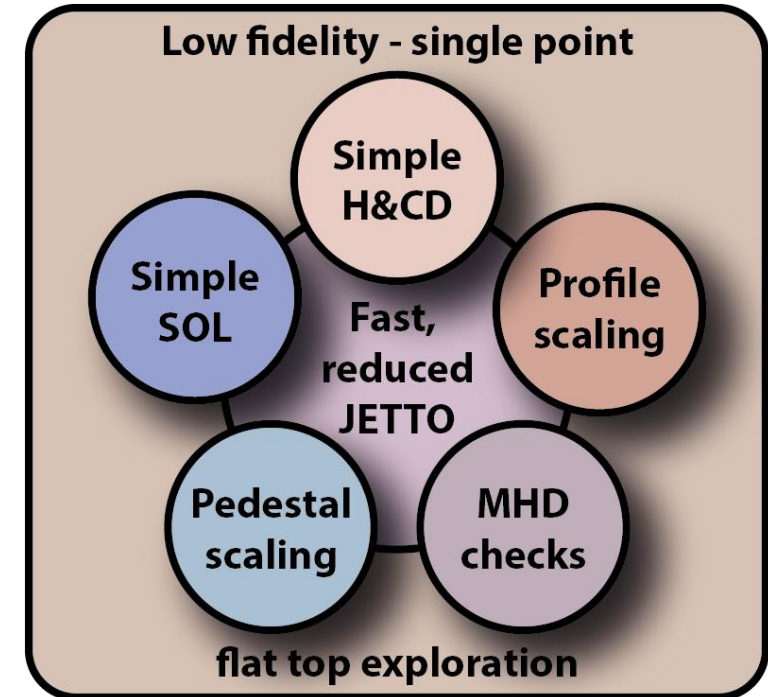
Tranche 3: Build
2032 - 2040

UK Atomic
Energy
Authority



1D transport code (JETTO) used as assumption integrator

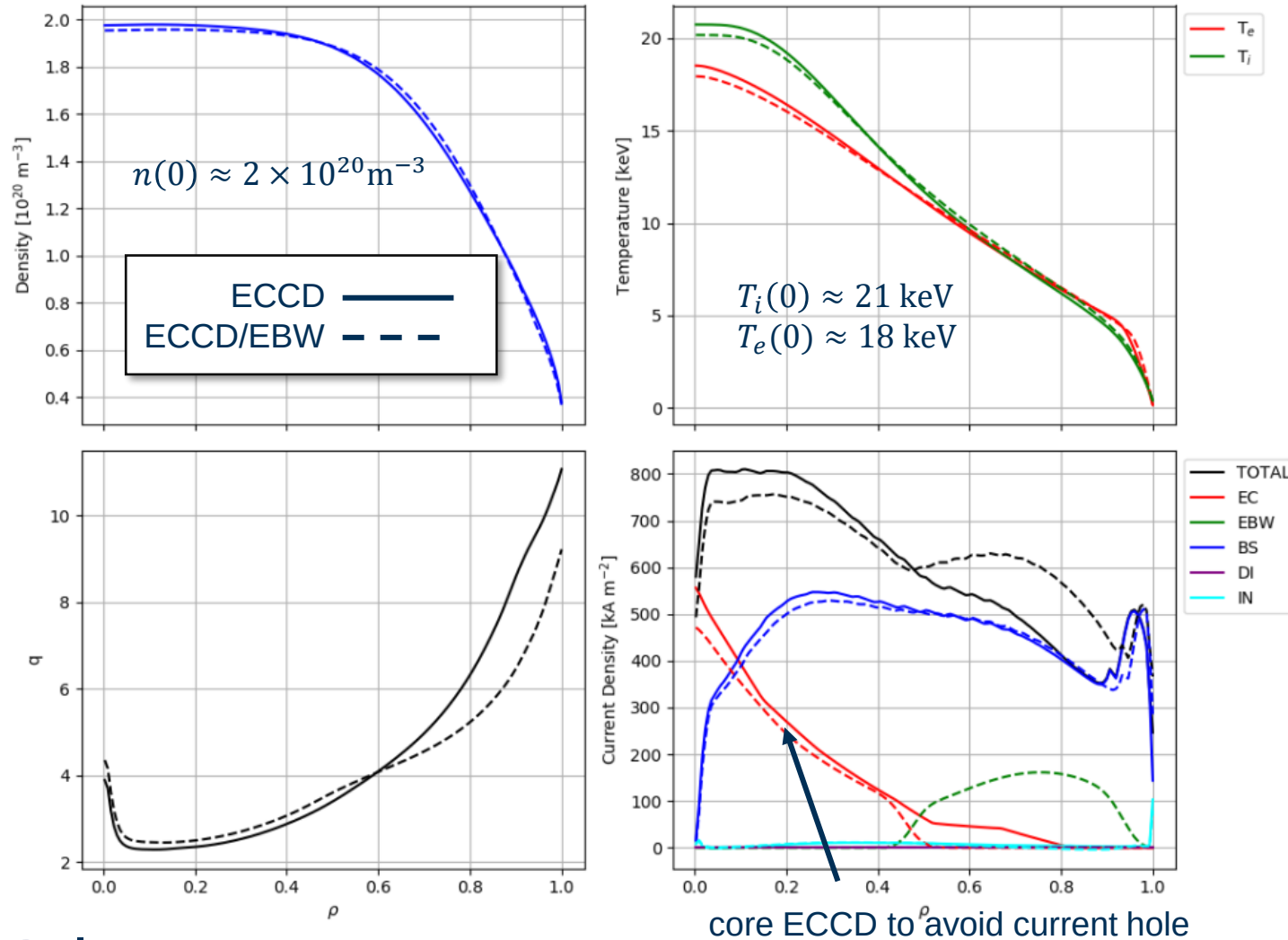
- STEP parameter regime outside validity of reduced transport models developed for present-day conventional tokamaks which don't capture electromagnetic (EM) turbulence expected to prevail in STEP
 - β_N is input
 - Other input parameters from systems code PROCESS
- Gyro-Bohm transport model produces profiles consistent with sources & sinks
 - Coefficients adapted to reflect dominant e^- transport as observed in present-day STs & suggested by gyrokinetic calculations
- Simplified heating & current drive models calibrated with higher fidelity calculations



Non-inductive flat-top operating points with access to ECCD & EBW $\Rightarrow B_T = 3.2$ T

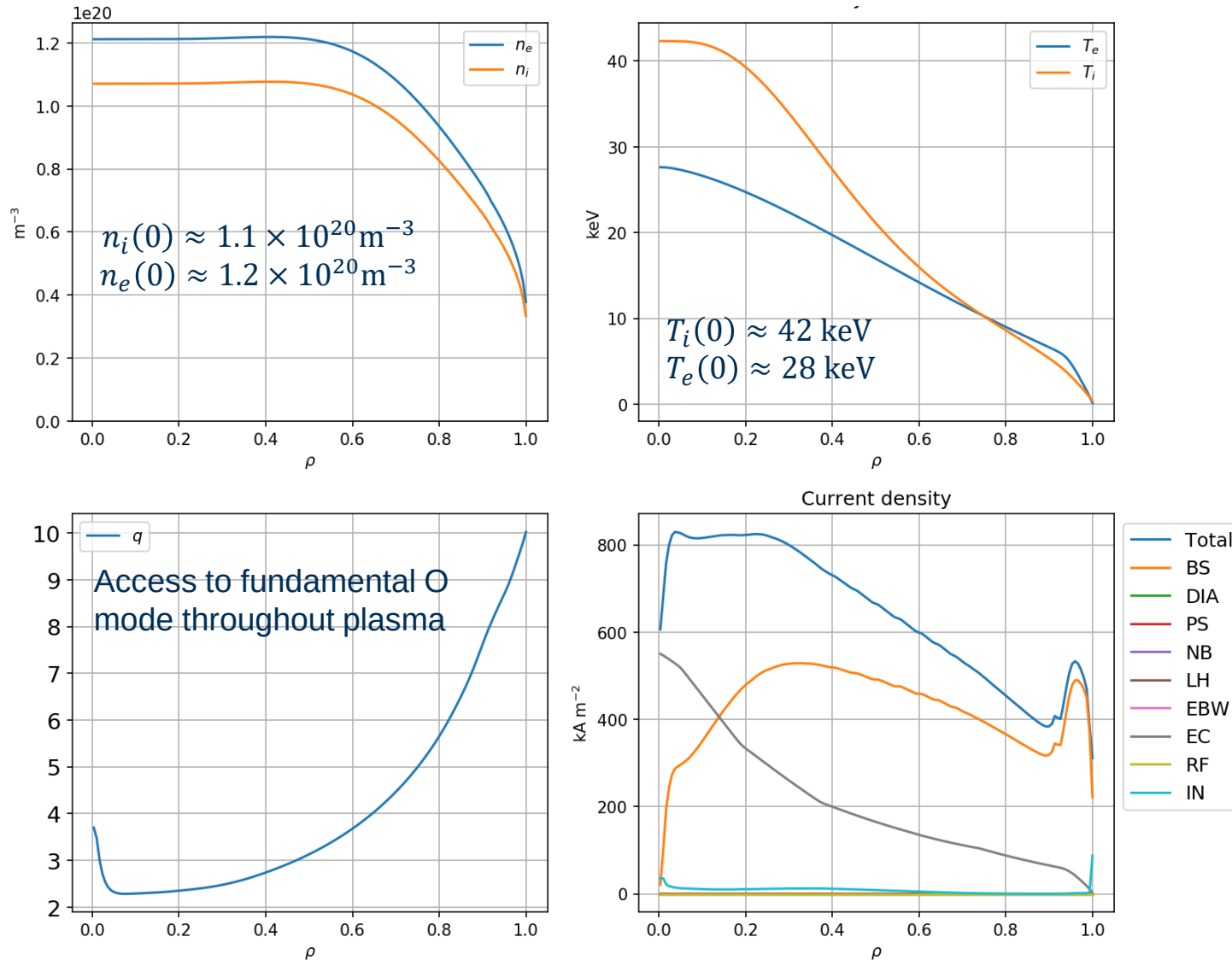
ECCD: Electron cyclotron heating & current drive

EBW: Electron Bernstein wave heating & current drive



H&CD	EC	EC/EBW
R_{geo} [m]	3.60	
A	1.8	
$B_T(R_{\text{geo}})$ [T]	3.2	
I_p [MA]	20.9	22.0
κ	2.93	
δ	0.59	0.50
P_{fus} [GW]	1.76	1.77
$P_{\text{net}}^{\text{el}}$ [MW]	188	182
P_{ECCD} [MW]	150	154
P_{rad} [MW]	338	341
Q	11.8	11.5
β_N	4.4	4.1
f_{BS}	0.88	0.78
$\bar{n}/n_{GW}$ [%]	100	94
$l_i(3)$	0.26	0.25
η_{CD}^{EC} [A/W]	0.016	0.027
η_{CD}^{EBW} [A/W]	N/A	0.034
$P_{\text{sep}}/R_{\text{geo}}$ [MW/m]	41	
$(H_{98} + H_{98}^*)/2$	1.35	1.19

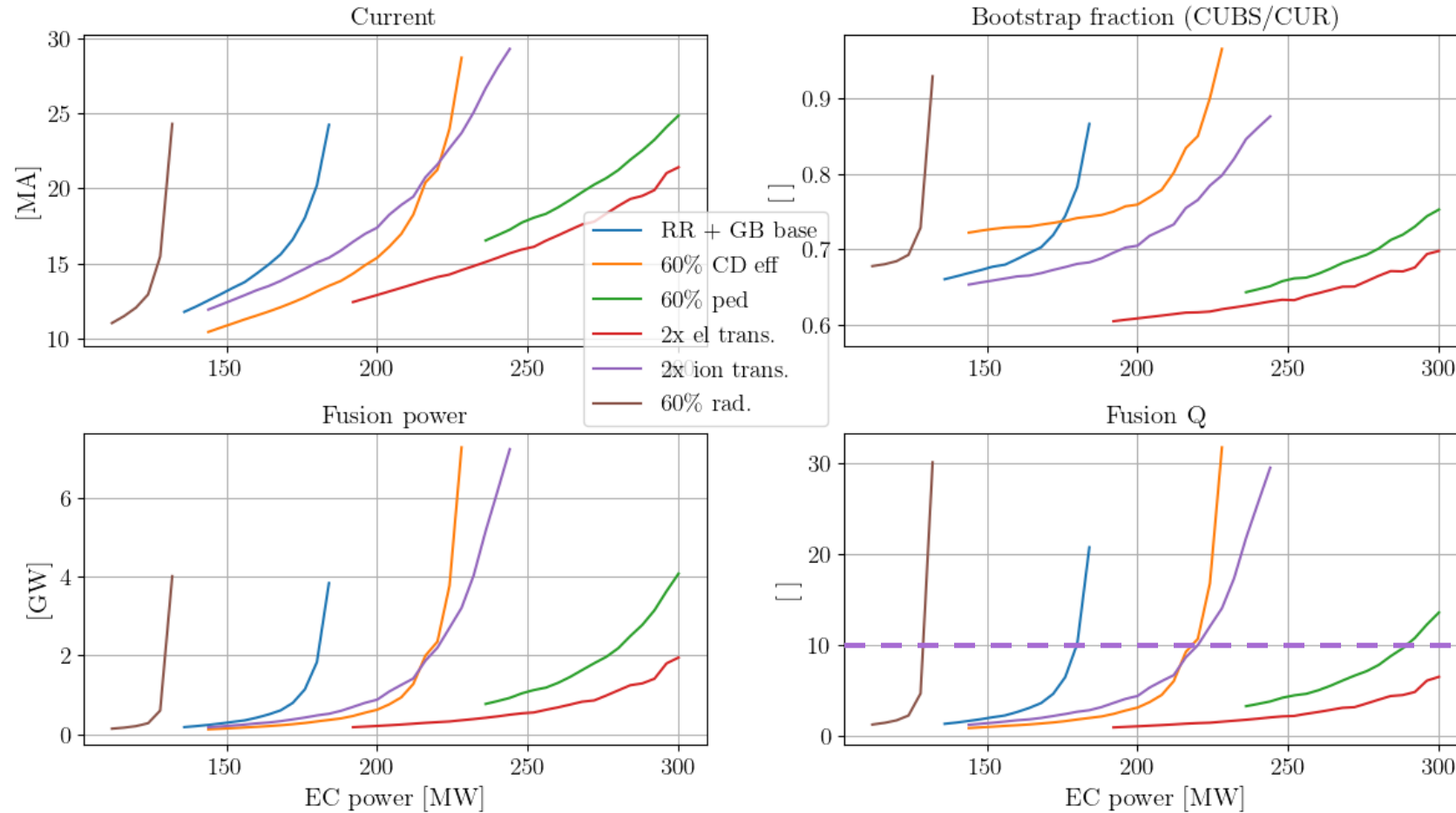
Low density hot ion flat-top operating point with access to fundamental ECCD



H&CD	EC
R_{geo} [m]	3.6
A	1.8
$B_T(R_{\text{geo}})$ [T]	3.2
I_p [MA]	22.8
P_{fus} [GW]	1.62
P_{aux} [MW]	150
P_{rad} [MW]	280
Q	10.1
β_N	4.6
$f_{BS} = I_{BS}/I_p$	0.76
$f_{GW} = \bar{n}/n_{GW}$	0.60
$l_i(3)$	0.30
η_{CD}^{EC} [A/W]	0.033
$P_{\text{sep}}/R_{\text{geo}}$ [MW/m]	41

Optimised using ITPA confinement scaling

“Predictive” uncertainty quantification shows margin in P_{aux} & P_{fus} needed to recover target Q



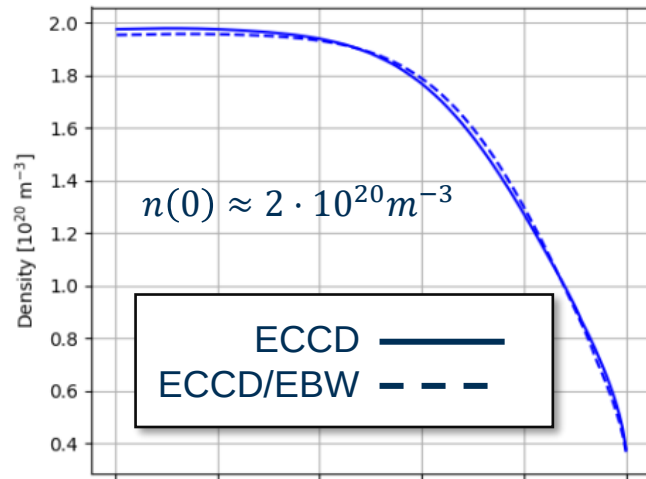
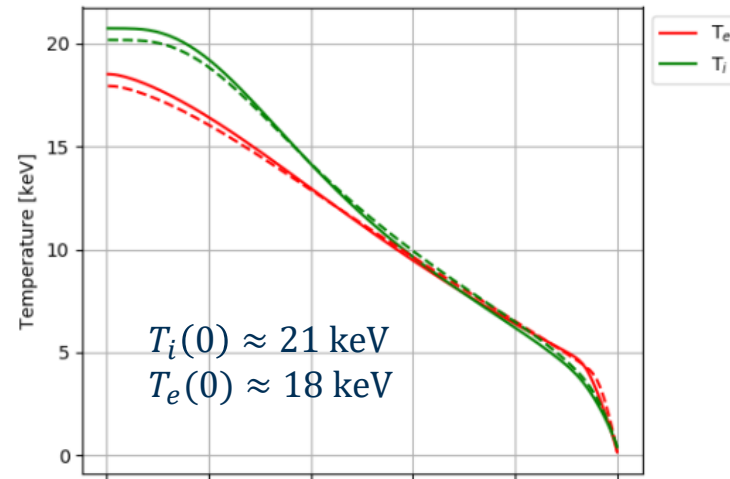
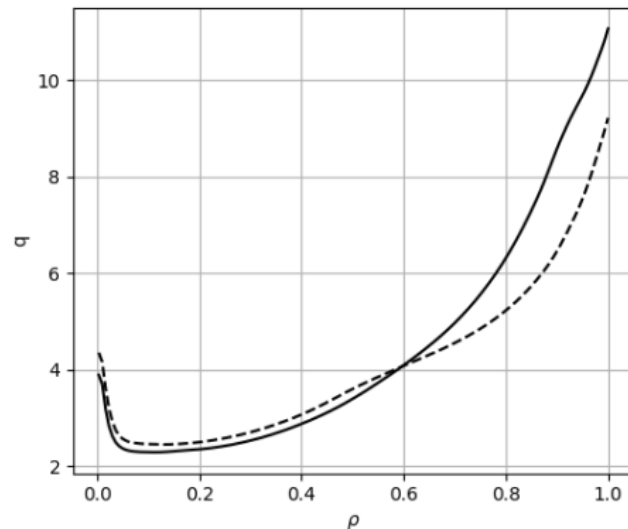
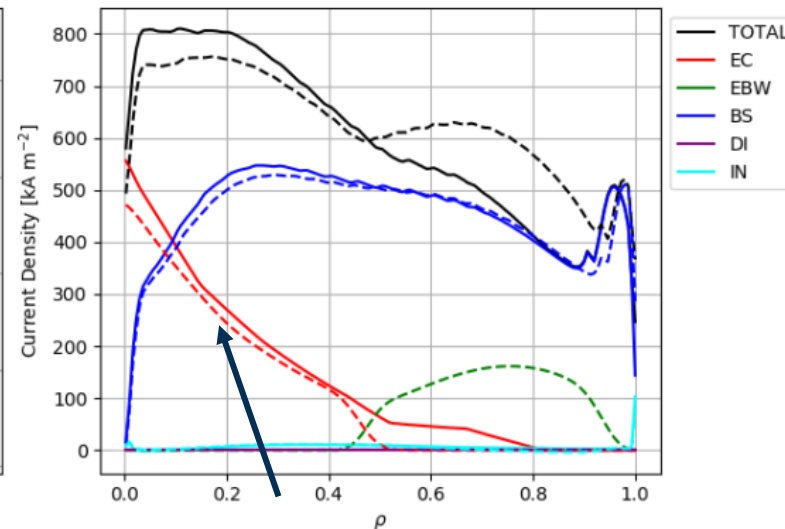
- Simple predictive setup:
 - Electrons: Rechester Rosenbluth (RR)
 - collisionless RR with $\delta B/B$ from [1] & connection length from [2]
 - Ions: gyro-Bohm
 - Predict T_i, T_e, n_e, I_p
 - Pedestal: $p_{ped} \propto I_p^{0.8} f_{GW}^{0.45}$
 $\Rightarrow P_{fus}$ very sensitive to I_p
- In base case $Q \approx 10$ needs $P_{aux} \sim 180$ MW & $P_{fus} \sim 2$ GW
- Performance sensitive to core radiation fraction

[1] J.F. Drake et al. PRL **44** (1980) 994
 [2] K. L. Wong et al. PRL **99** (2007) 135003

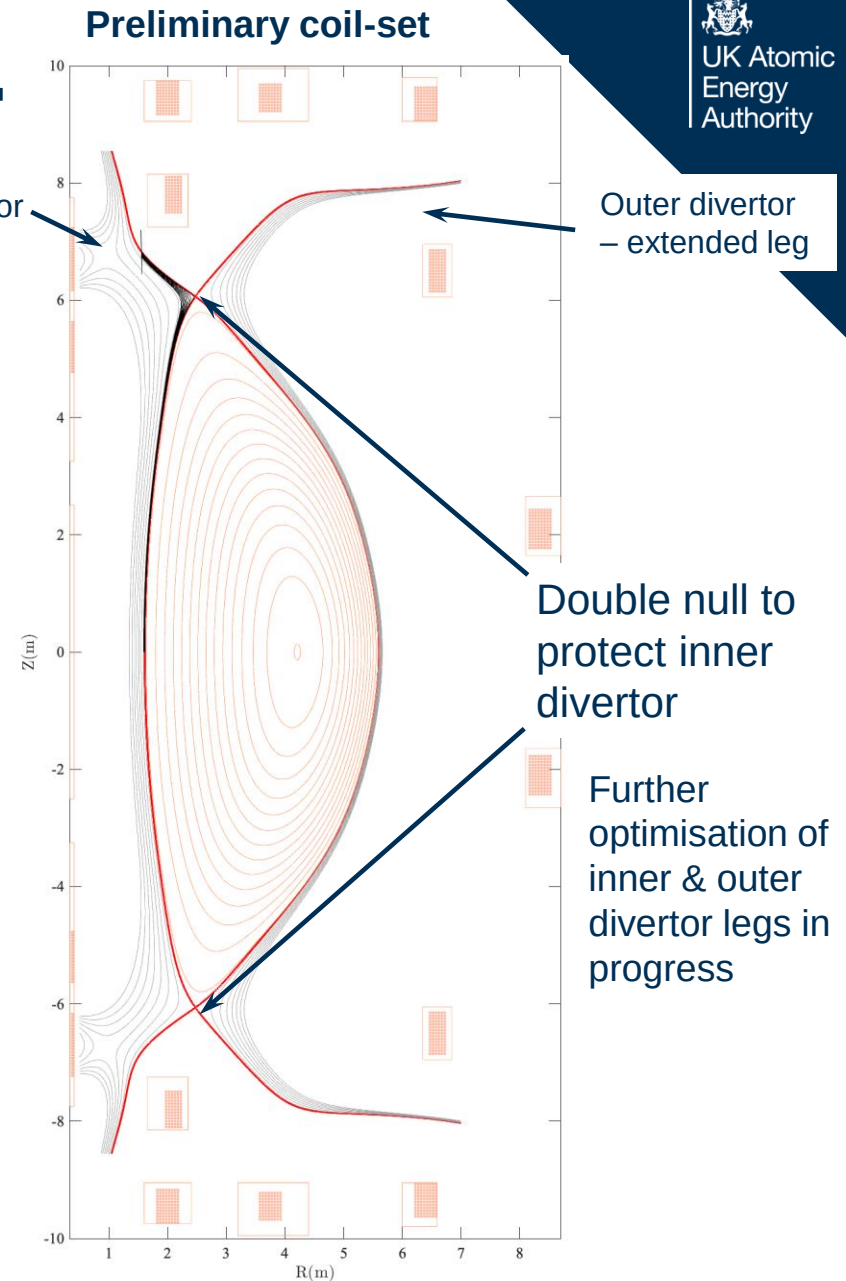
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EBW: Electron Bernstein wave heating and current drive

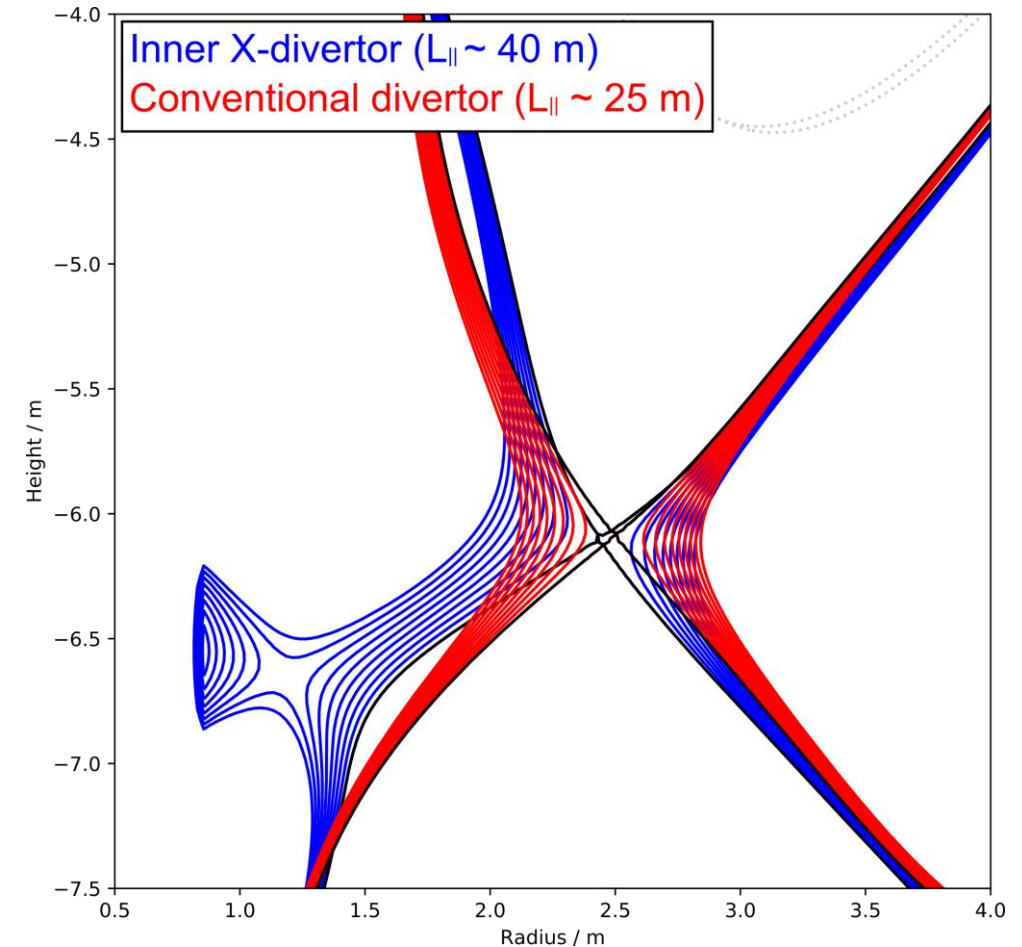


core ECCD to avoid current hole



X-divertor may provide optimum exhaust solution on inboard side

- SOLPS-ITER code used to assess level of detachment & target heat loads in various divertor geometries
- Inboard X-divertor poses engineering challenge but has significant advantages:
 - when strike point radius $\sim 1.7\text{m}$, connection length $L_{\parallel}$ is nearly doubled
 - offers best performance with fewest compromises in terms of neutral trapping & peak heat loading
 - for details, see talk by S. Henderson, *Physics Drivers of the STEP divertor concept design*, this workshop



S. Henderson, A. Hudoba

Active control of Resistive Wall Modes (RWMs) allows access to higher β_N

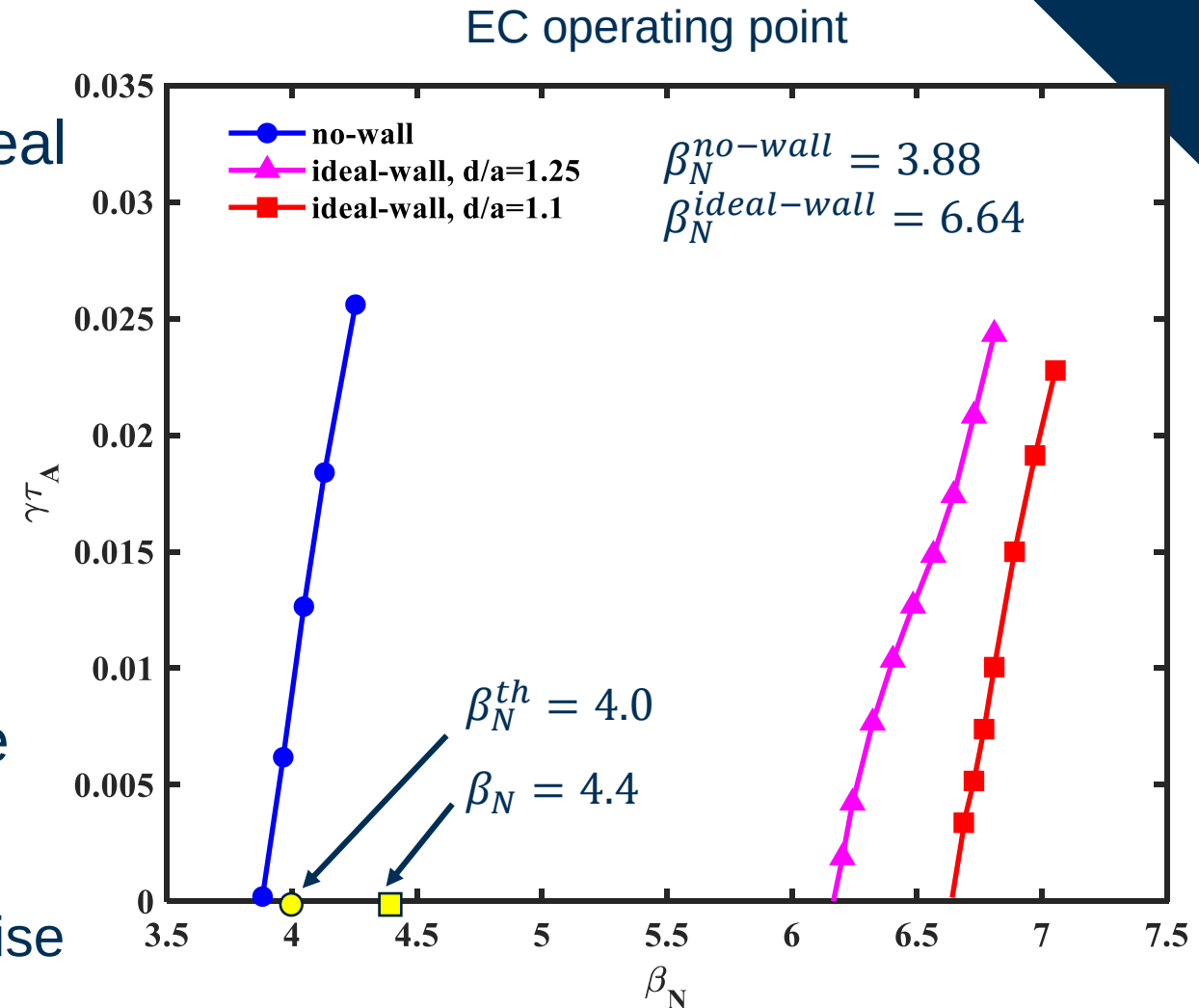
- STEP scenarios are marginally above ideal stability limit without conducting wall:

$$\beta_N > \beta_N^{no-wall}$$

- Conducting wall reduces growth rate to resistive values $\Rightarrow$ controllable:

$$\beta_N < \beta_N^{ideal-wall}$$

- Active in-vessel coils can keep amplitude of instability well below disruption limit
 - 6 mid-plane picture-frame coils
 - Control system modelled with system noise



Active control of Resistive Wall Modes (RWMs) allows access to higher β_N

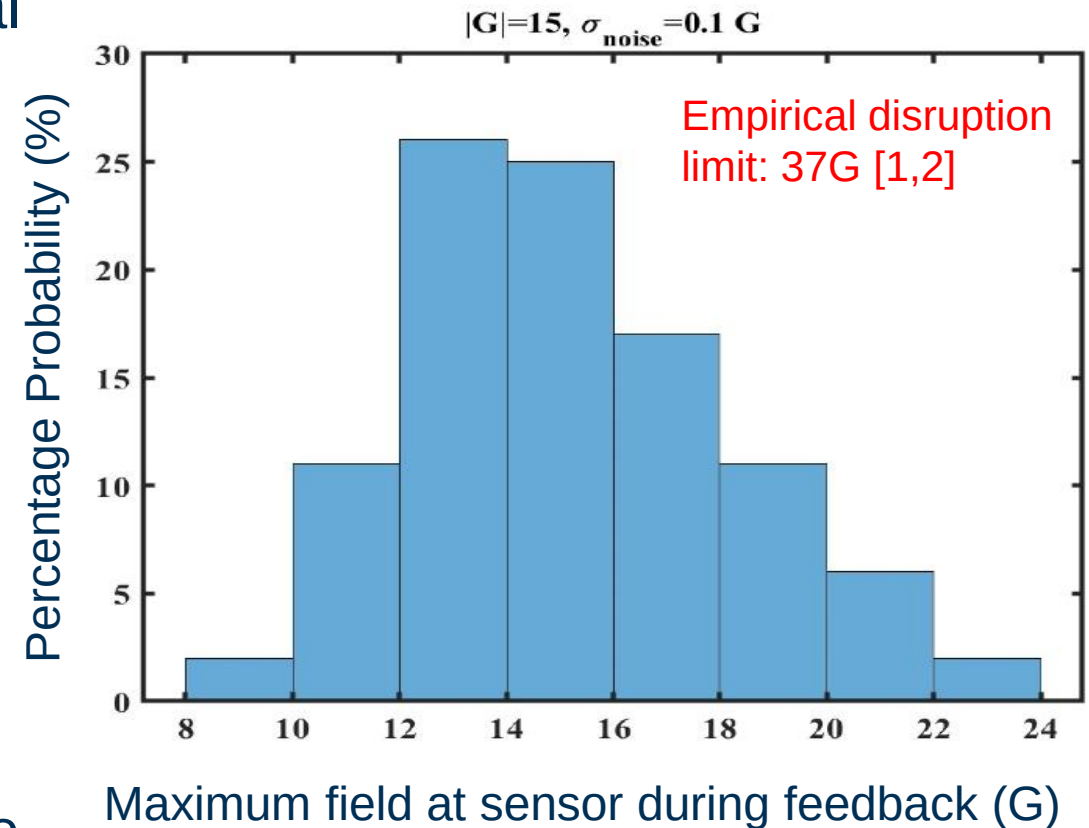
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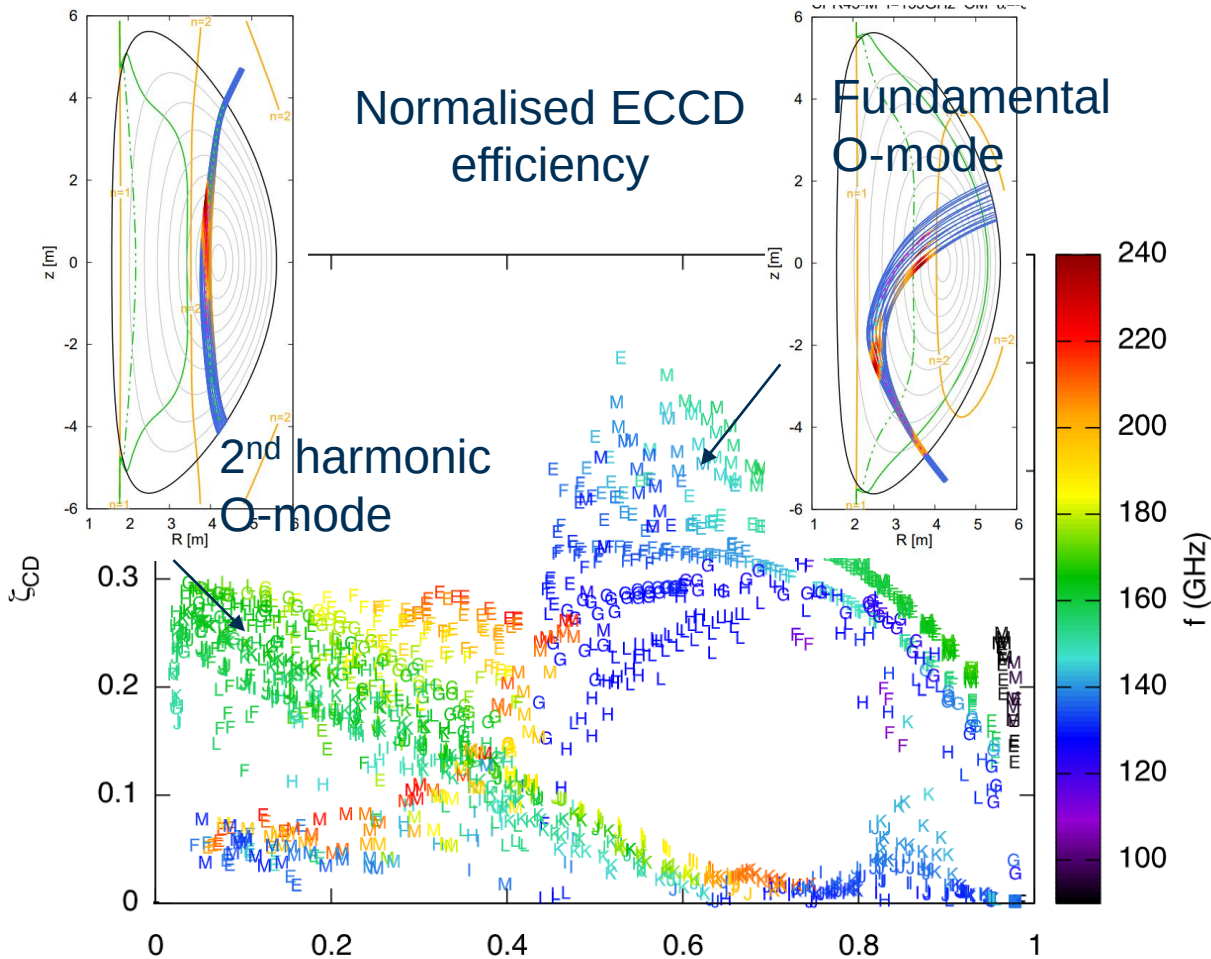
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[1] P.C. de Vries et al. Nucl. Fusion **56** (2016) 026007

[2] G. Xia et al. Proc. 48th EPS Plasma Phys. Conf. (2022)

ECED O-mode covers full radius; EBW only covers $\rho > 0.4$ but at 3 - 4 $\times$ higher normalised efficiency



$$\zeta_{CD} = 32.7 \frac{\eta[A/W] n_e [10^{20} m^{-3}] R [m]}{T_e [keV]}$$

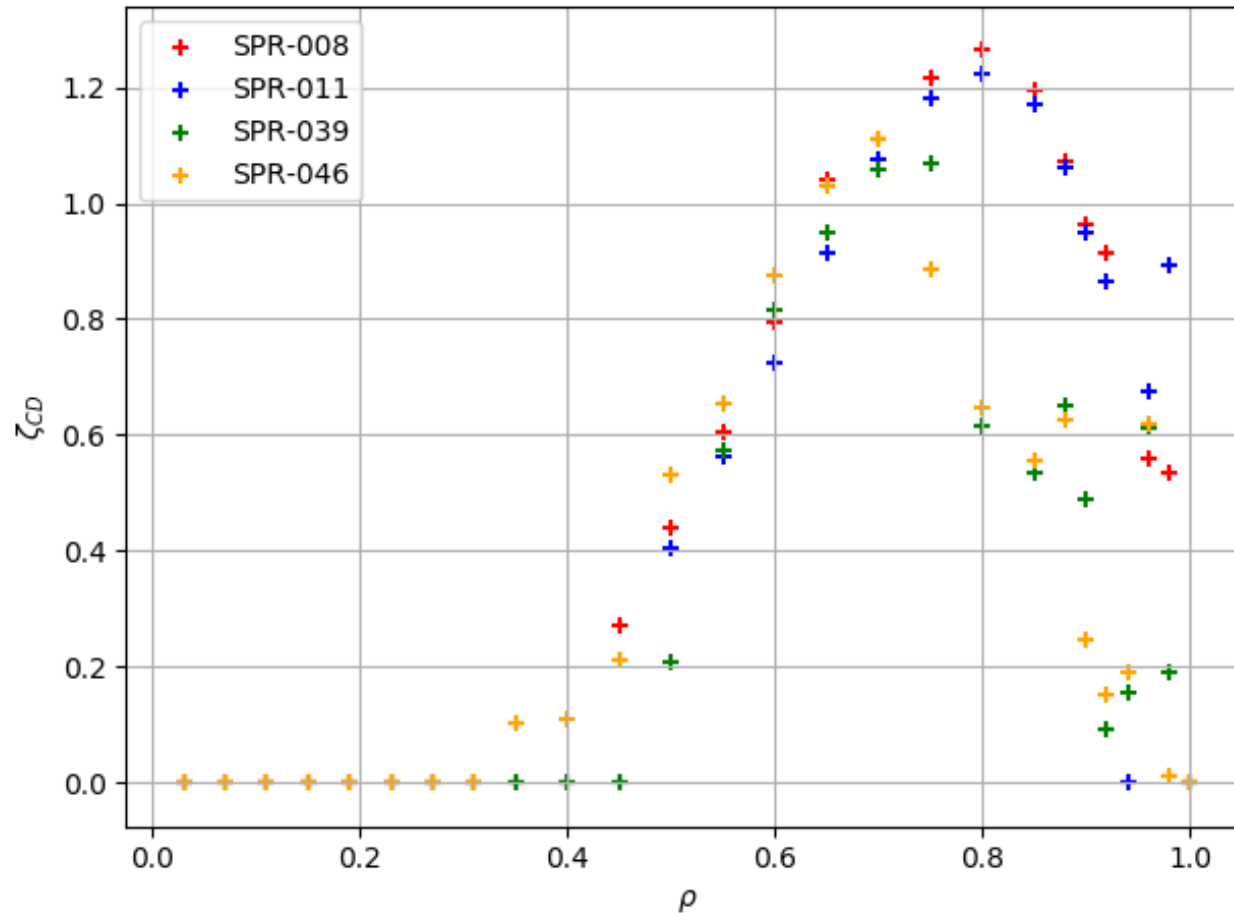
- ECED: scan with GRAY code for multiple launch positions
 - Low-field side O-mode launch from above/below midplane allows access through magnetic field well
 - High-field side absorption negates particle trapping degradation for off-axis current drive
- EBW: full wave calculation using GENRAY + CQL3D
 - High central T_e makes $\rho < 0.4$ inaccessible
 - 2nd harmonic with dominant Ohkawa current drive [1]

ρ_{tor} L. Figini (CNR), M. Henderson, S. Freethy

[1] G. Taylor et al. Phys. Plasmas **11** (2004) 4733

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Normalised EBW CD efficiency



$$\zeta_{CD} = 32.7 \frac{\eta[A/W] n_e [10^{20} m^{-3}] R[m]}{T_e [keV]}$$

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D. Speirs (U. Strathclyde), T. Wilson, M. Henderson, S. Freethy

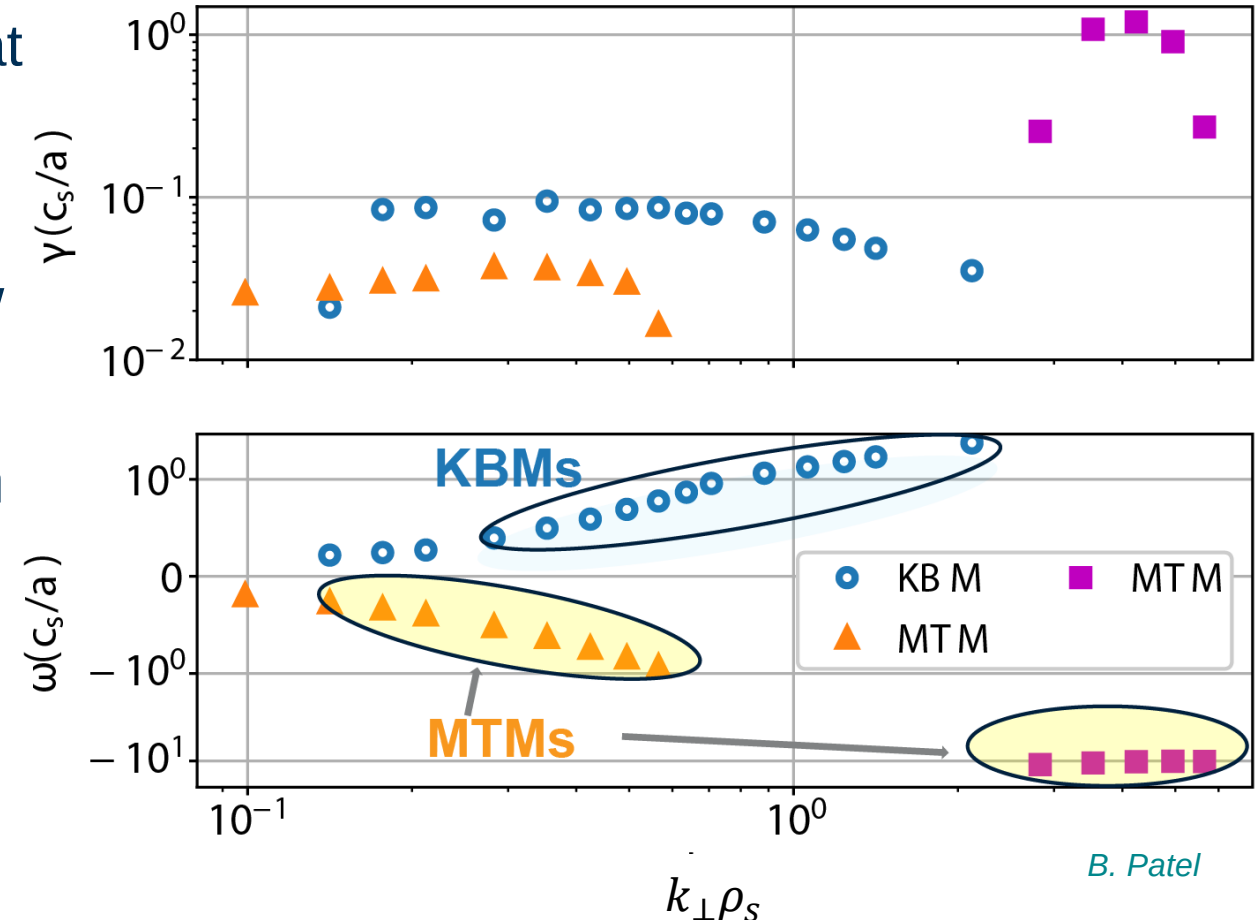
[1] G. Taylor et al. Phys. Plasmas **11** (2004) 4733

Microstability in STEP differs significantly from JET, DEMO & ITER

B.S. Patel et al. Nucl. Fusion **62** (2022) 016009

- Electromagnetic turbulence dominates
 - 2 types of micro tearing mode (MTM) at low $k_{\perp}\rho_s$ (i-MTM) & high $k_{\perp}\rho_s$ (e-MTM)
 - kinetic ballooning modes (KBMs)
- KBMs have highest growth rates at low $k_{\perp}\rho_s$ but may be stabilised by flow & β'
- MTMs mainly drive conductive electron heat transport via turbulent radial reconnection of field lines (magnetic flutter [1])

[1] J.D. Callen PRL **39** (1977) 1540



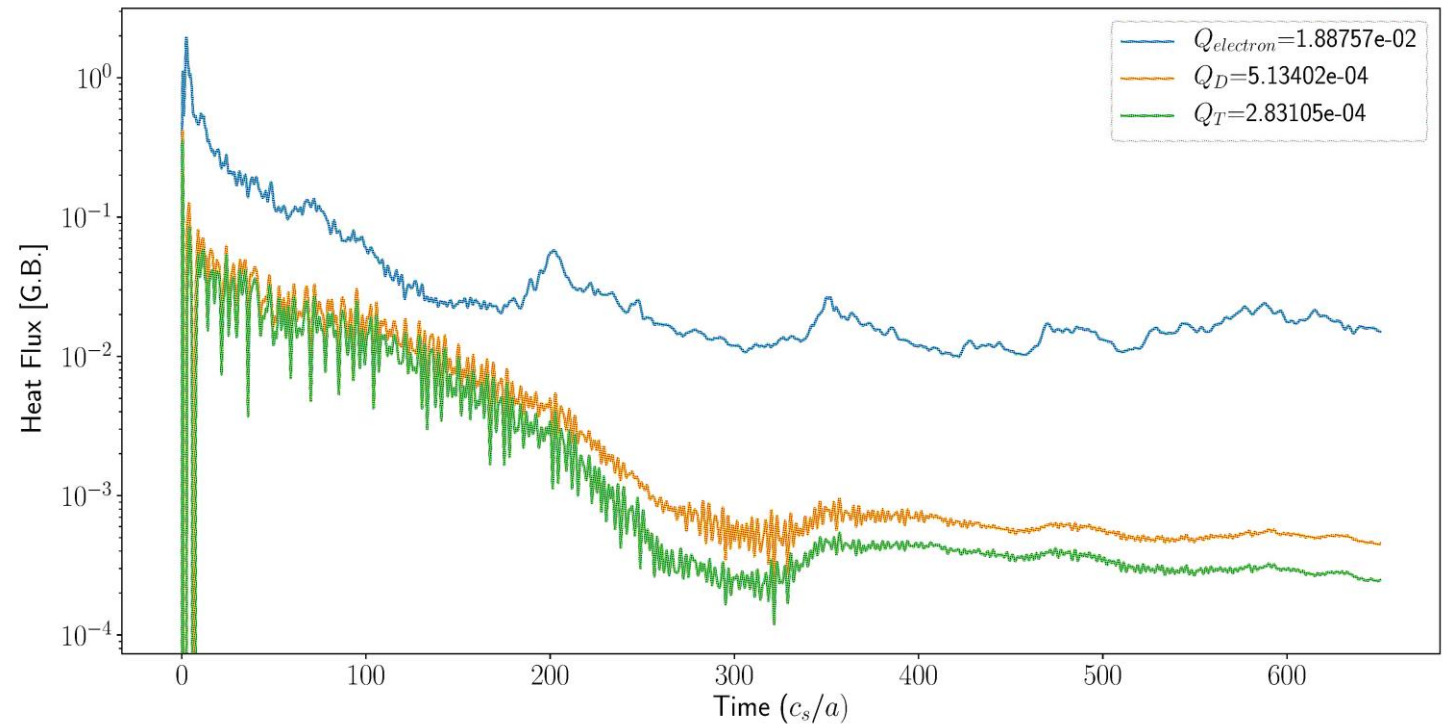
B. Patel

Recent non-linear local gyrokinetic simulations that exclude KBMs show saturation at $Q_e \sim 0.01 Q_{gB}$

- Earlier local nonlinear electromagnetic gyrokinetic simulations for STEP predicted very large magnetic flutter transport from low wavenumber MTMs

D. Dickinson, Varenna Fusion Theory Conference, 2022

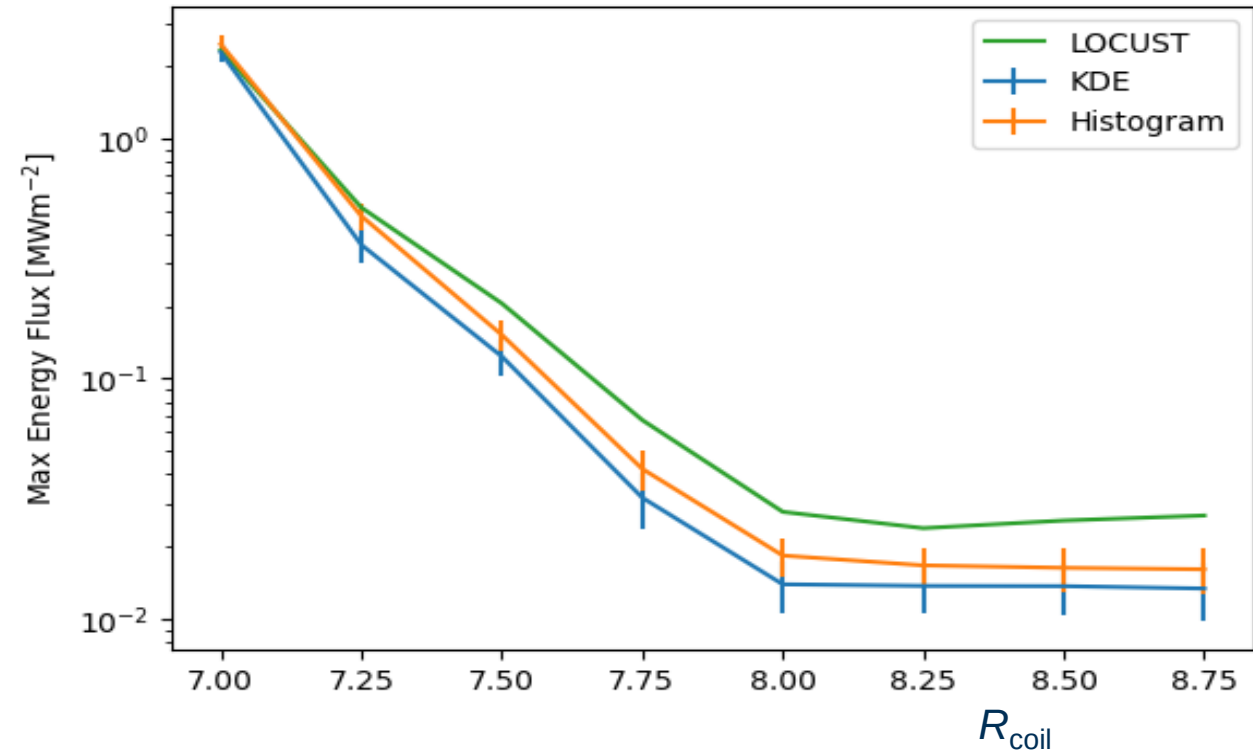
- Recent investigations that exclude KBMs reveal strong sensitivity of heat flux to parallel dissipation
 - GENE simulations with $B_{||} = 0$ (no KBMs) show greatly reduced Q_e & $Q_{D,T} \ll Q_e$
- i-MTMs & KBMs have different radial scales
 - vastly increases computational cost of nonlinear simulations that include both types of mode



D. Kennedy, B. Patel, C. Roach, D. Dickinson (U. York), M. Giacomini (U. York)

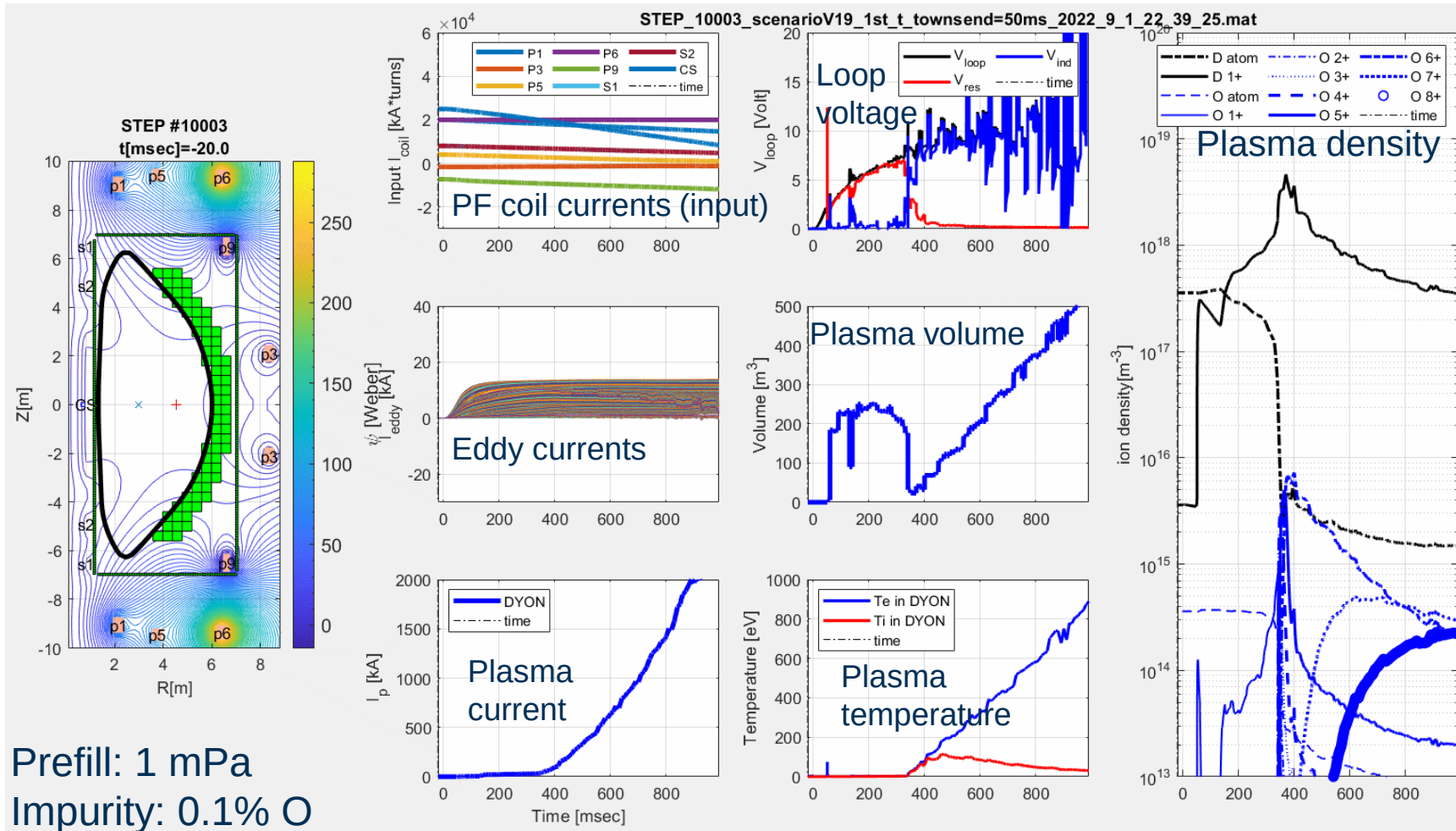
Power loads due to α -particle losses used to constrain parameters of TF & ELM control coils

- Full α -particle orbits in 3D fields tracked using LOCUST code
- TF ripple-induced α -particle losses & distribution of associated power loads on 1st wall calculated for $N = 16$ picture-frame coils with range of outer limb radii R_{coil}
- Maximum loads occur on low field side main chamber wall which can tolerate up to $\sim 1 \text{ MWm}^{-2}$ in total (including EM radiation)
- Power loads acceptably low for $R_{\text{coil}} \geq 8.0 \text{ m}$
- Coils designed for RWM control & error field correction may also be used for ELM control
- α -particle loss & power load calculations for 3D fields needed for ELM control coils underway – significant uncertainty in coil configurations required for suppression of Type I ELMs



A. Prokopyszyn, Steady-state alpha-particle flux on the walls of the Spherical Tokamak for Energy Production, this workshop

Burn-through achieved after 350 ms with $V_{loop} \sim 6.5$ V using hexapole null at $R = 2.5$ m



- Small ~ 9 Vs solenoid used for initiation & target plasma production
- DYON code used to simulate burn-through
 - supported by free boundary equilibrium solver (FIESTA)
 - self-consistent calculation of eddy currents & Townsend break down
- ECRH pre-ionisation & heating expected to lower V_{loop} requirement

H.T. Kim

STEP plasma work improving confidence in feasibility of ST-based fusion power plant

Fast concept turn-around has made it possible to explore variety of whole plant concepts using 1D transport code & integrated modelling tools

Three flat-top operating points have been defined, trading confinement risk against heating & current drive maturity

Comprehensive tool set for scenario modelling has helped to reduce risks in exhaust, stability, current drive, core plasma & α -particle confinement

Other STEP presentations at this workshop (Wednesday)

Stuart Henderson – Physics drivers of the STEP divertor concept design (oral)

Emmi Tholerus – Analysing flat-top operational space of STEP reactor design (oral)

Alex Prokopyszyn - Steady-state alpha-particle flux on the walls of the spherical tokamak for energy production (oral)

Alex Fil – Disruption runaway electrons generation & mitigation attempts in STEP preferred concept (poster)

Žiga Štancar – Modelling of the STEP ramp-down phase with JETTO (poster)

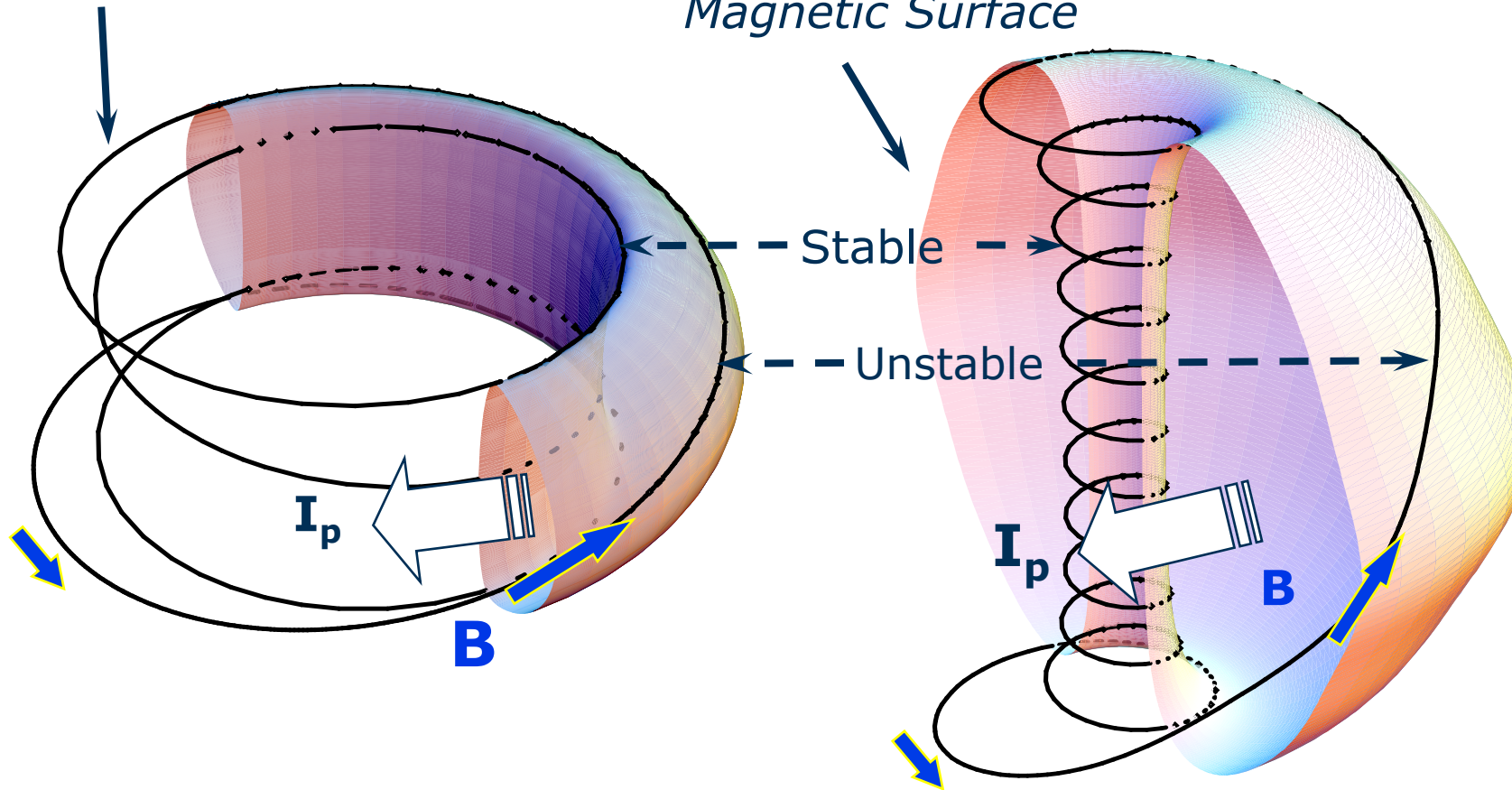
Sam Bakes - Free boundary equilibrium modelling of disruption scenarios for STEP (poster)

Back-up slides

Why choose ST route to fusion power plant?

Magnetic Field Line

Magnetic Surface



Conventional Tokamak

Spherical Tokamak

cartoon courtesy of Y.M. Peng

- STs allow higher plasma pressure in lower toroidal field $\Rightarrow$ higher β
- STs operate naturally at higher elongation κ

$$P_{fus} \propto n^2 \langle \sigma v \rangle_{DT} \propto p^2 \propto \beta_T^2 B_T^4$$

$$\text{with } \beta_T \sim A^{-1/2} (1 + \kappa^2) \beta_N^2 / f_{BS}$$

J.E. Menard et al. NF **37** (1997) 595

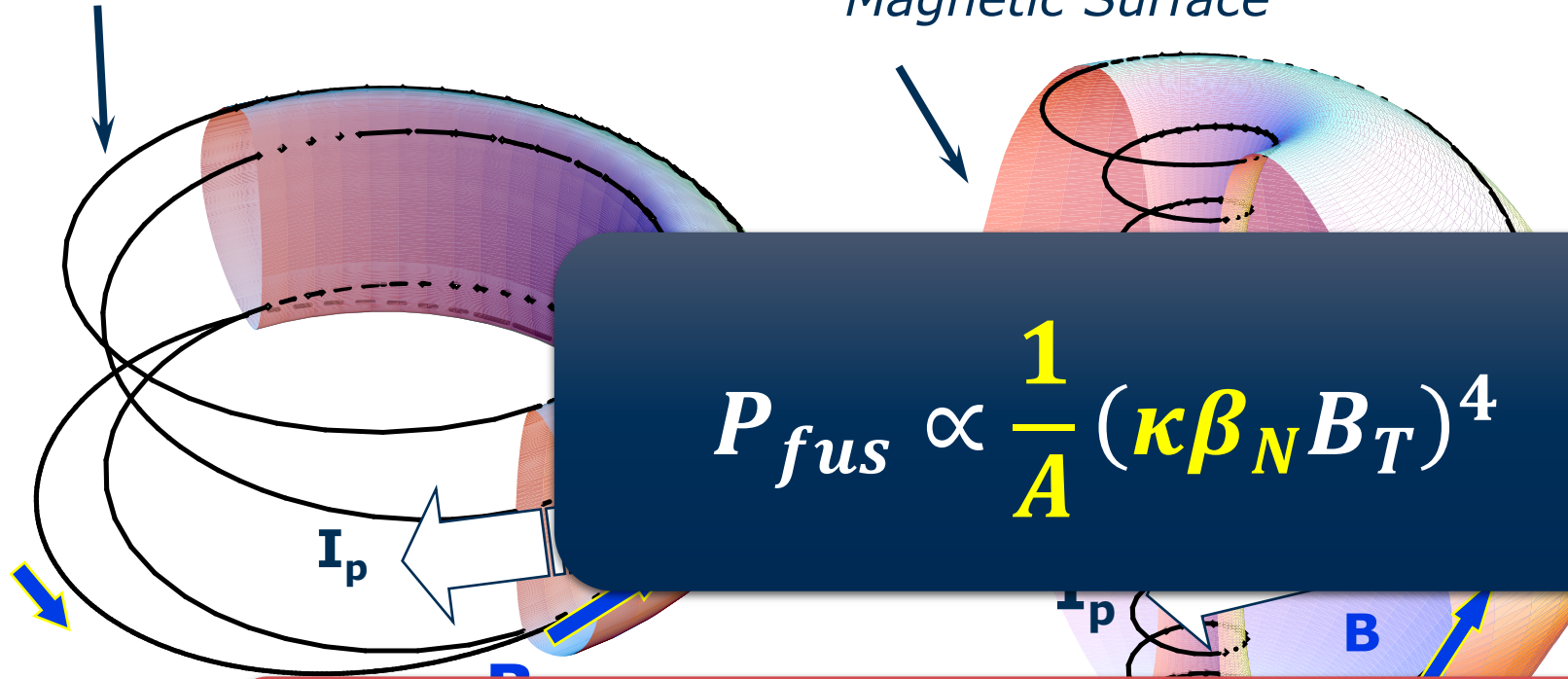
$$\beta_N = \frac{\beta}{I_p / a B_T} \quad f_{BS} = I_{BS} / I_p$$

(I_{BS} : bootstrap current)

Why choose ST route to fusion power plant?

Magnetic Field Line

Magnetic Surface



$$P_{fus} \propto \frac{1}{A} (\kappa \beta_N B_T)^4$$

- STs allow higher plasma pressure in lower toroidal field $\Rightarrow$ higher β
- STs operate naturally at higher elongation κ

$$P_{fus} \propto n^2 \langle \sigma v \rangle_{DT} \propto p^2 \propto \beta_T^2 B_T^4$$

Operation at high β_N and high κ require broad current profile (low l_i) plus robust vertical & resistive wall mode (RWM) control

$$(1 + \kappa^2) \beta_N^2 / f_{BS}$$

NF 37 (1997) 595

Conventional Tokamak

Spherical Tokamak

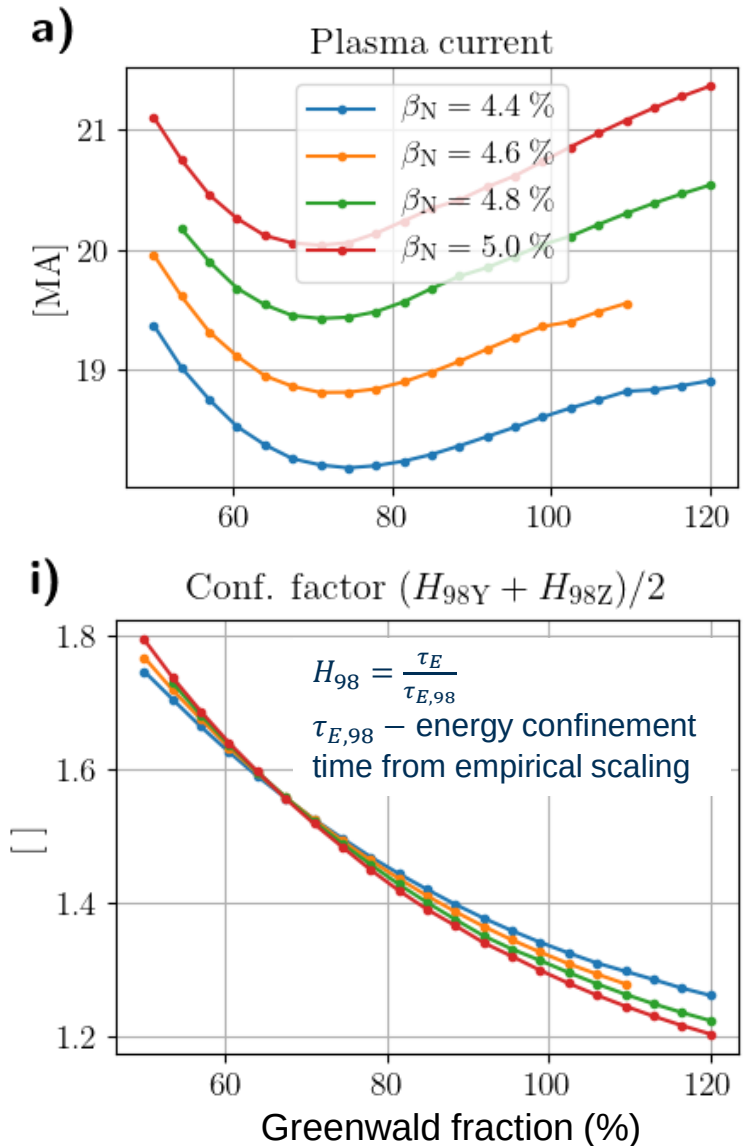
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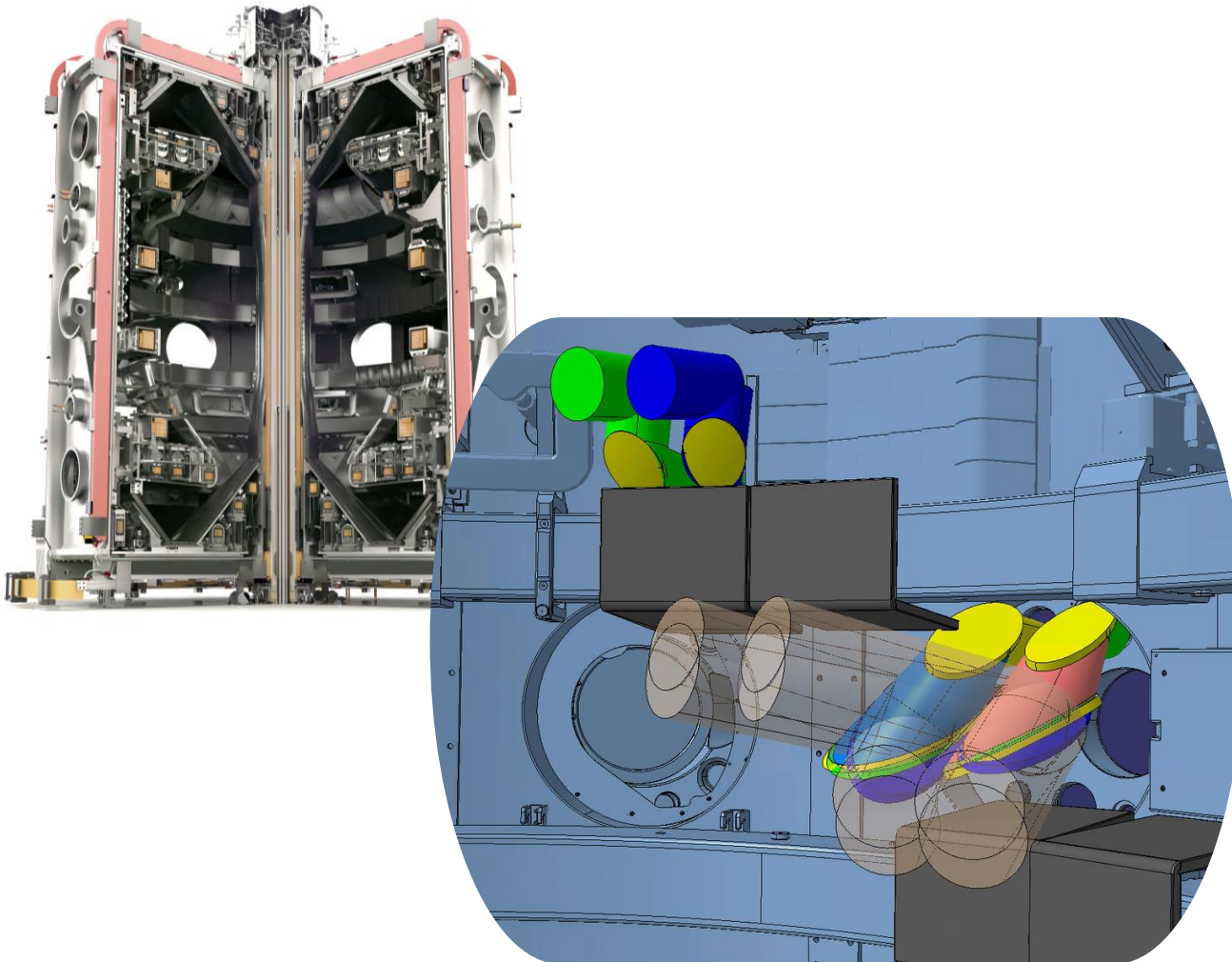
cartoon courtesy of Y.M. Peng

Confinement uncertainty traded against uncertainty in current drive efficiency

- Minimising requirement on plasma confinement at fixed β_N requires operation at high density $\Rightarrow$ high Greenwald fraction f_{GW}
- Scans performed at fixed P_{aux} with fuelling adjusted to achieve target f_{GW}
- Electron cyclotron current drive (ECCD) efficiency η_{CD} is low at high density $\Rightarrow$ requires high bootstrap current fraction f_{BS}
- High f_{BS} requires high confinement $H_{98} \Rightarrow$ no operational point with pure ECCD & $1 \lesssim H_{98} \lesssim 1.2$
 - Lower plot: non-intuitive scaling of $H_{98} \sim \beta_N^{0.14}$
- Electron Bernstein waves (EBWs) expected to provide higher η_{CD}
 - wider operational space but moves risk to current drive technique with less mature physics basis



1.8 MW microwave system will be installed on MAST Upgrade to test EBW performance on ST



- To address open physics questions, two gyrotrons manufactured by Kyoto Fusioneering will be installed in UKAEA MAST Upgrade spherical tokamak:
 - Dual frequency (28 / 34.8 GHz)
 - Up to ~900kW & ~4.5s pulse length
 - In-vessel launcher capable of:
 - on- / off- axis current drive
 - co- / counter- current drive
 - electron heating (no current drive)
 - solenoid-free start-up
- System will be ready for operation in 2024

M. Henderson, S. Freethy

ST-specific challenges for plasma

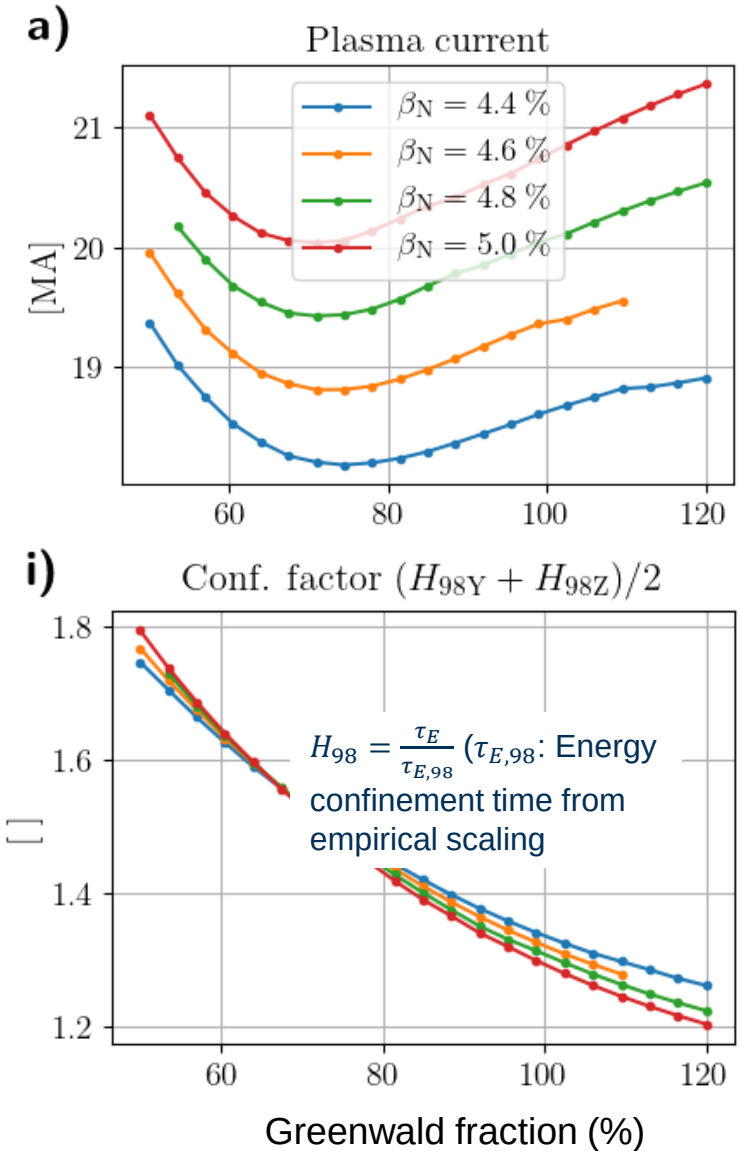
- Compact design $\Rightarrow$
 - Less surface for tritium breeding
 - Less surface for handling heat & particle fluxes $\Rightarrow$ **alternative divertor design**
 - Less space for solenoid $\Rightarrow$ **Need to design for largely non-inductive pulse**
 - Non inductive current drive: challenge that all tokamak reactors need to face eventually!
- Optimisation of pressure-driven current fraction (bootstrap) f_{BS} to reduce need for auxiliary current drive
 $\Rightarrow$ **operation at high normalised pressure β_N & elongation κ**
 - Operation at high $\beta_N \Rightarrow$
 - Need to operate at elevated $q_{min} > 2 \Rightarrow$ efficient **off-axis current drive**
 - access to 2nd stability, avoidance of low n/m neoclassical tearing modes (NTM)
 - Need to actively control resistive wall modes (RWM)
 - Operation at high elongation $\Rightarrow$
 - Need to operate at low internal inductance $l_i \Rightarrow$ efficient **off axis current drive**

ST-specific challenges for the Plasma

- Compact design $\Rightarrow$
 - Less surface for tritium breeding
 - Less surface for handling heat & particle fluxes $\Rightarrow$ **alternative divertor design**
 - Less space for **neutron pulse**
 - Non inductive operation $\Rightarrow$ **face eventually!**
- Optimisation of auxiliary current $\Rightarrow$ **operation at**
 - Operation without large edge localised modes (type-I ELMs) **reduce need for**
 - Operation & control of self organised plasma at high radiation & bootstrap fraction
 - Control with scarce actuators & sensors
 - Need to operate at high temperature (NTM)
 - Need to actively control resistive wall modes (RWM)
- Operation at high elongation $\Rightarrow$
 - Need to operate at low internal inductance $l_i \Rightarrow$ efficient **off axis current drive**

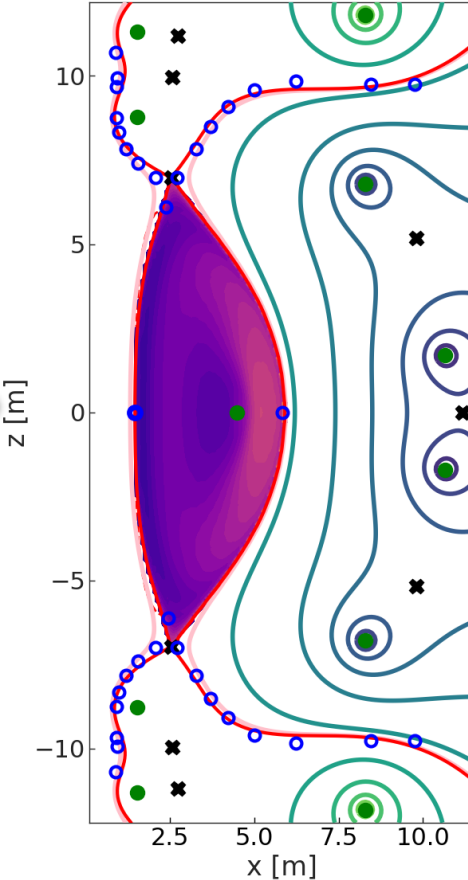
Trade uncertainty on confinement against uncertainty of current drive method

- To minimise the requirement on plasma confinement at fixed β_N requires operation at high density $\Rightarrow$ high Greenwald fraction f_{GW}
- Electron Cyclotron current drive (ECCD) efficiency, η_{CD} , is low at high density $\Rightarrow$ high bootstrap current fraction f_{BS}
- High f_{BS} requires high confinement assumption $H_{98} \Rightarrow$ no operational point with low $1 \lesssim H_{98} \lesssim 1.2$ with ECCD
 - At high f_{BS} and $Q \Rightarrow$ non intuitive scaling of $H_{98} \sim \beta_N^{0.14}$
- Electron Bernstein Wave expected to have much higher $\eta_{CD} \Rightarrow$ wider operational space but moves risk to current drive technique with less mature physics basis



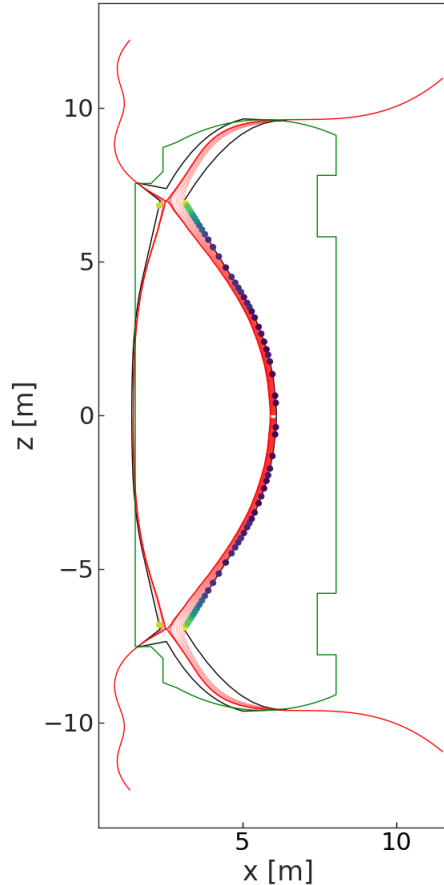
Development of an engineering concept optimiser

PF coil sets optimisation: A. Hudoba,
S. Bakes, G. Voss, S. Kahn



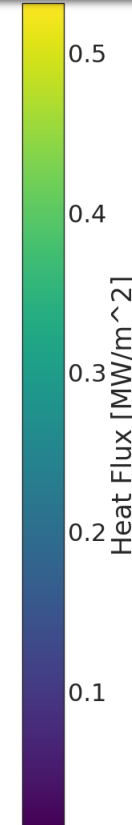
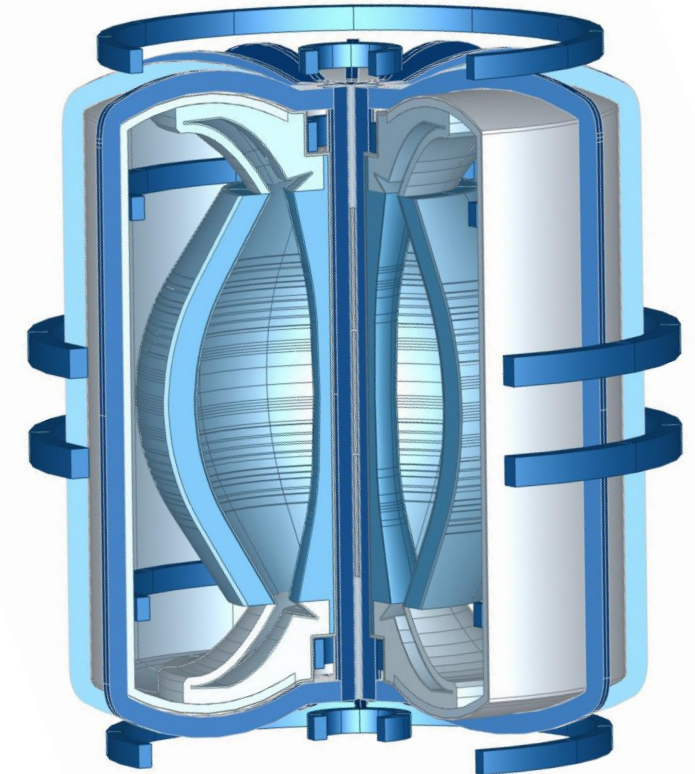
Free boundary
Equilibrium

Wall: D. Vaccaro, J. Harrison



First wall

PF coil optimisation in progress
TF optimisation with selected shapes
reducing TF ripple



Reduced 3D
CAD

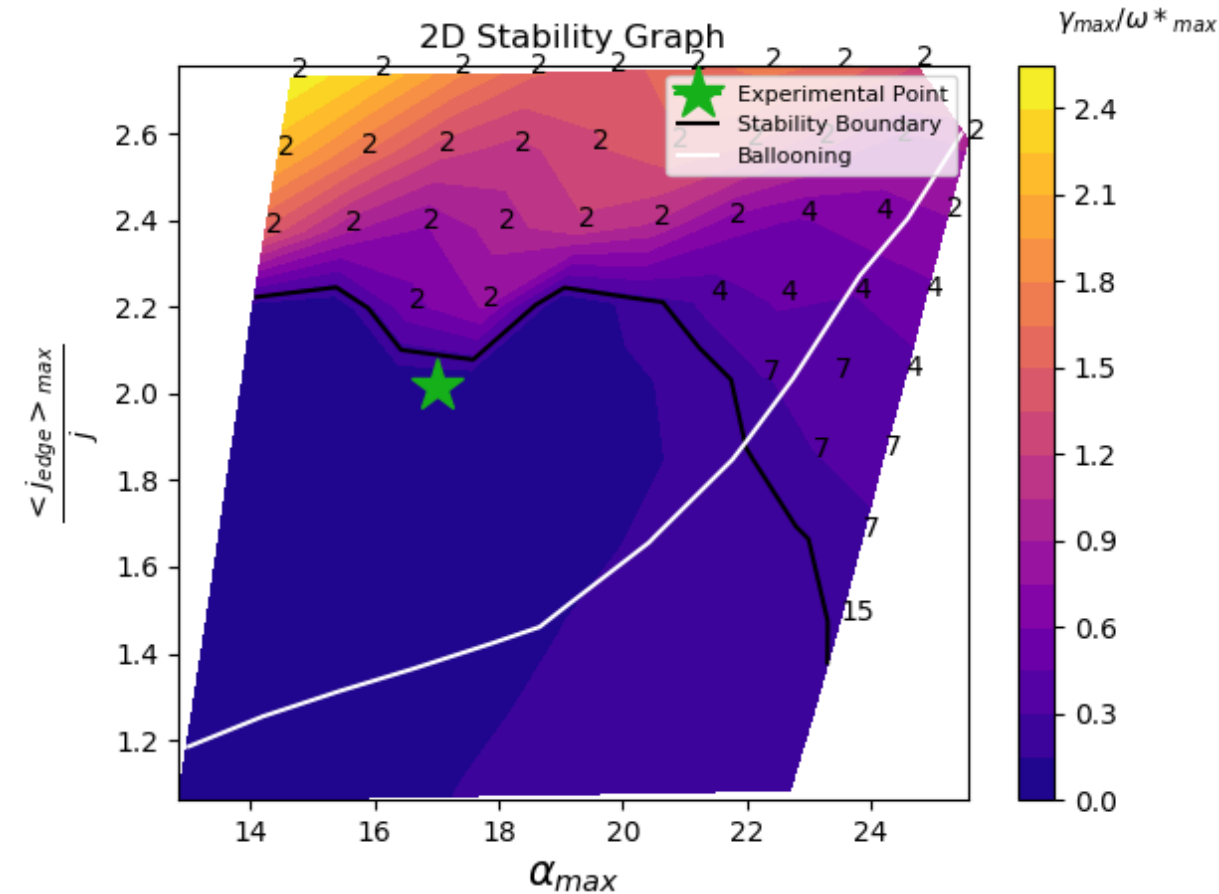
PROCESS
JETTO

STEP will have a small $\psi \sim 9 \text{ Vs}$ solenoid to allow break down & initial limited ramp-up

- DYON code used to assess break-down
 - So far pessimistic assumption with no ECRH
 - Field-null evolution included
- First assessment $V_{loop} \sim 20 \text{ V}$ and $\psi \sim 1.4 \text{ Vs} \Rightarrow I_p \sim 0.4 \text{ MA}$
 - ECRH/ECCD assist should make it possible to reduce V_{loop}
- Inductive limiter phase to generate initial target plasma.
 - Modelling shows $I_p \sim 1 \text{ MA}$ with $\psi \sim 2 \text{ Vs}$
 - Need to increase plasma volume & elongation to allow divertor formation for non-inductive phase.

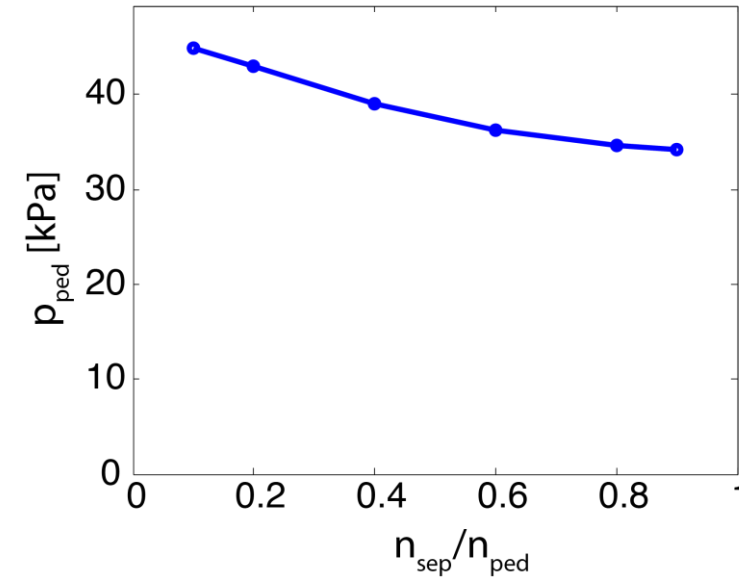
STEP plasmas fulfil necessary condition to allow access to quiescent H-mode

- Large edge localised modes (type I ELMs) cannot be tolerated in STEP $\Rightarrow$ **ELM free regimes needed**
 - Energy fluence per ELM much higher than melt limit of W
- Better pedestal performance relieves requirements on core performance.
 - Pedestal pressures that are too high may lead to intolerable α -particle losses
- Pedestal performance found to be insensitive to separatrix density n_{sep} $\Rightarrow$ **main driver of plasma triangularity δ**



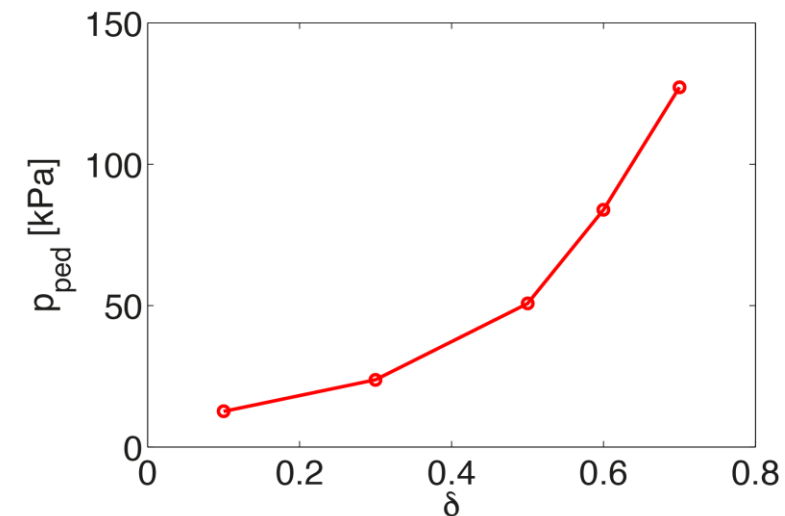
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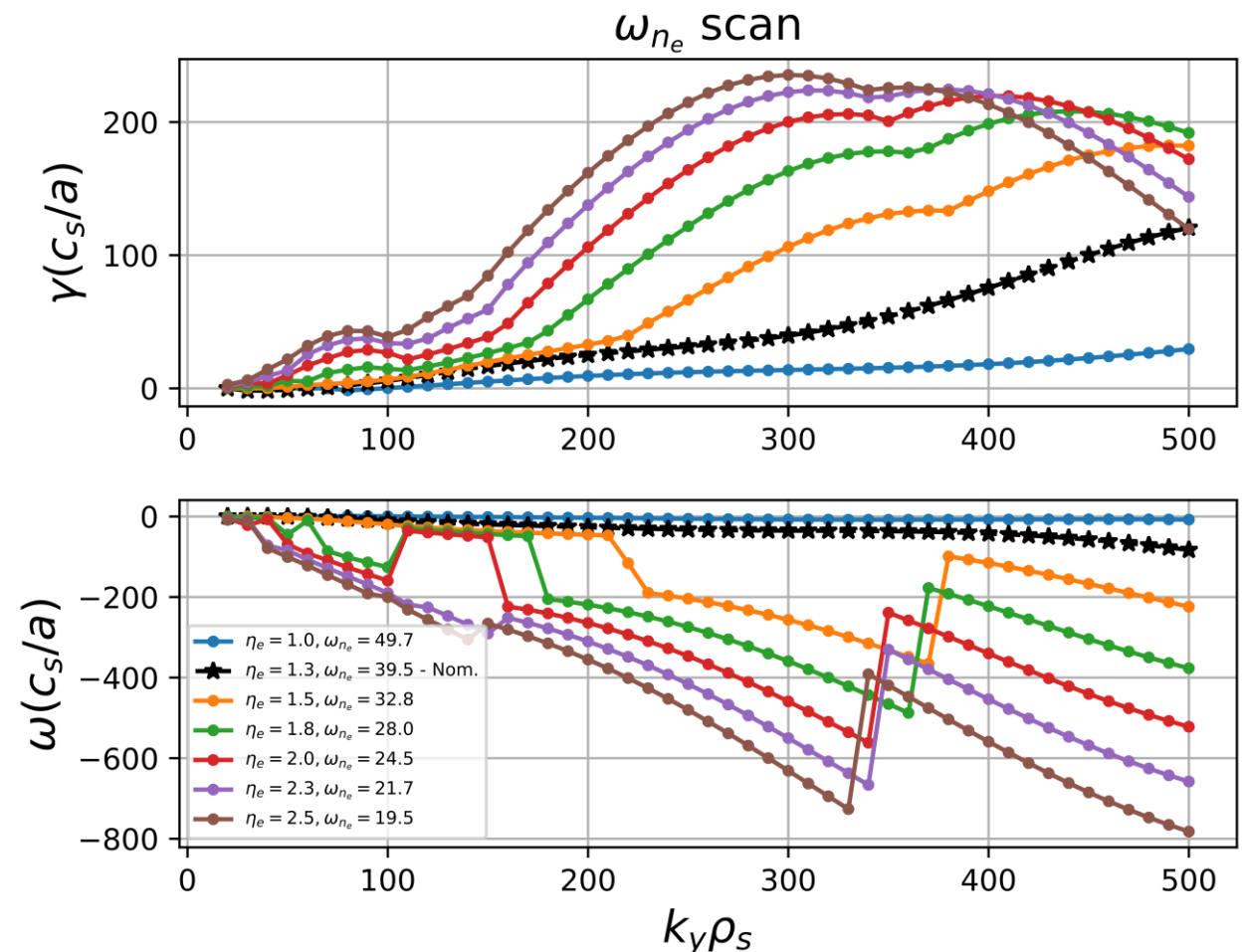
$$R_{geo} = 2.5 \text{ m}$$

$$B_t = 2.3 \text{ T}$$



Linear gyrokinetic analysis suggests that ETGs could limit pedestal gradients

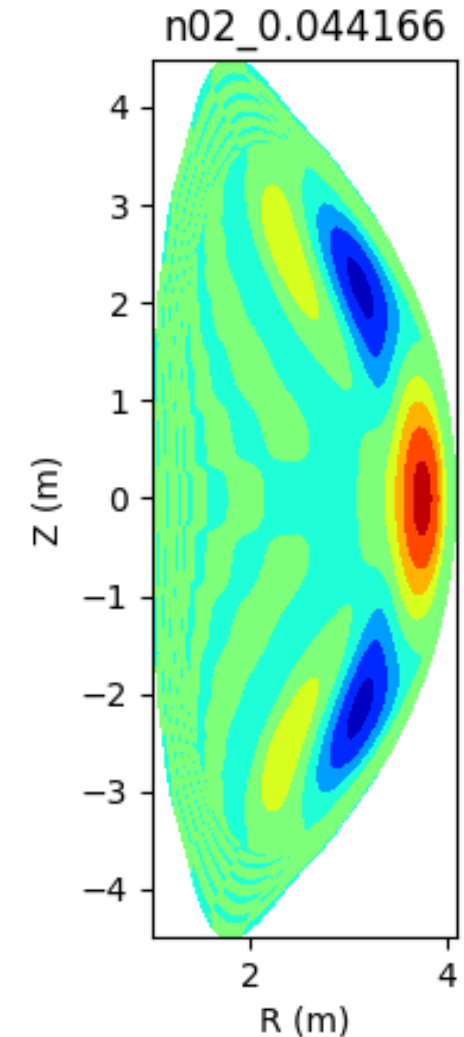
- First linear GK analysis for typical STEP pedestal using gyrokinetic code GENE
 - Motivated by findings from JET
- Profiles close to marginal stability for slab-ETGs
 - Increasing $\eta_e = L_{n_e}/L_{T_e}$ by varying density gradient reveals robust slab-ETG modes



ETG: Electrostatic electron temperature gradient driven mode

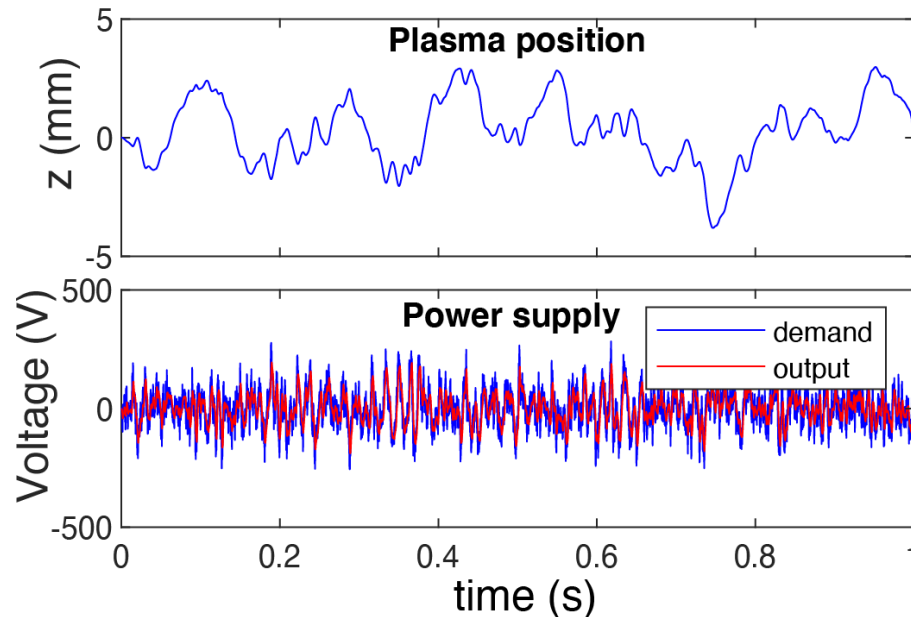
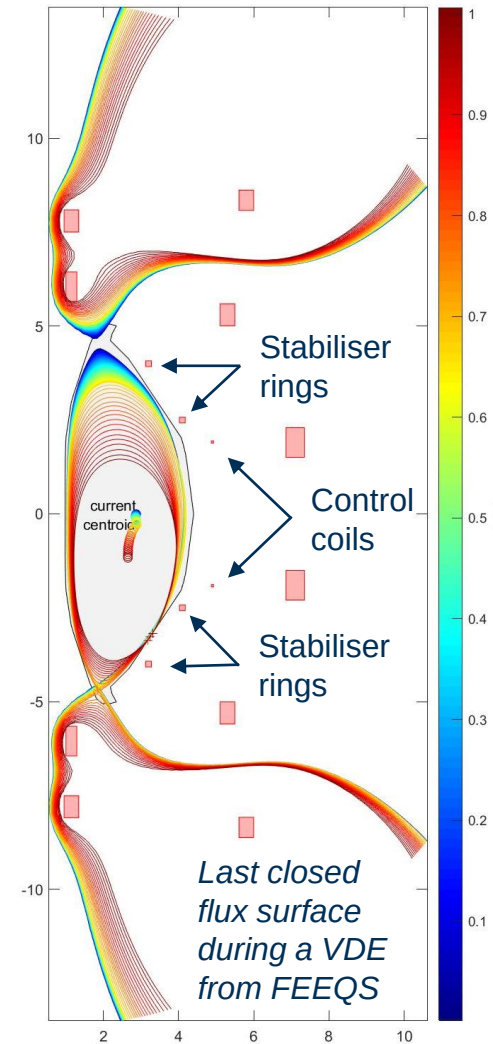
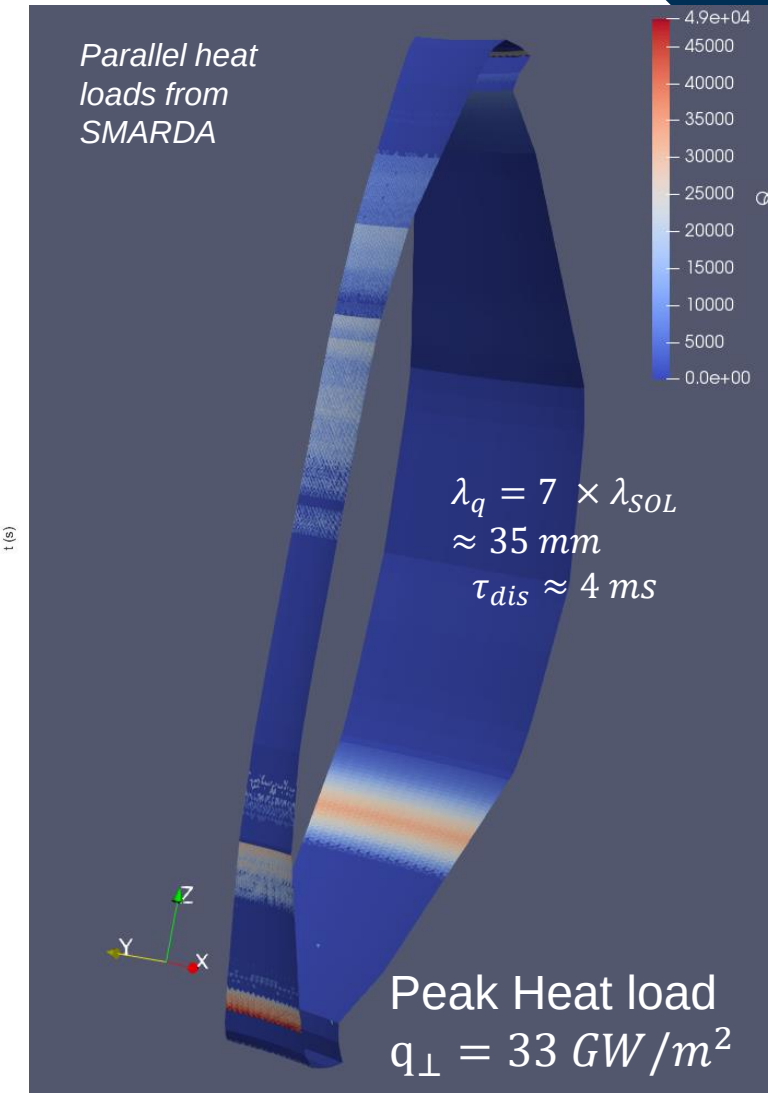
TAE thermal damping stronger than growth

- STEP equilibria have wide spectrum of toroidal Alfvén Eigenmodes (TAES), some with broad radial structure & relatively high α -particle drive $\frac{\gamma_{TAE}}{\omega} \sim 0.2$
 - TAEs can degrade α -particle confinement
- However TAEs predicted to be stable in flat-top due to high β & consequent strong thermal ion damping
 - May become problematic for phase when current is ramped-up & thermal ion β is low
 - modes could be driven by hot electrons
 - Higher frequency (ellipticity/triangularity induced) Alfvén eigenmodes also need to be investigated



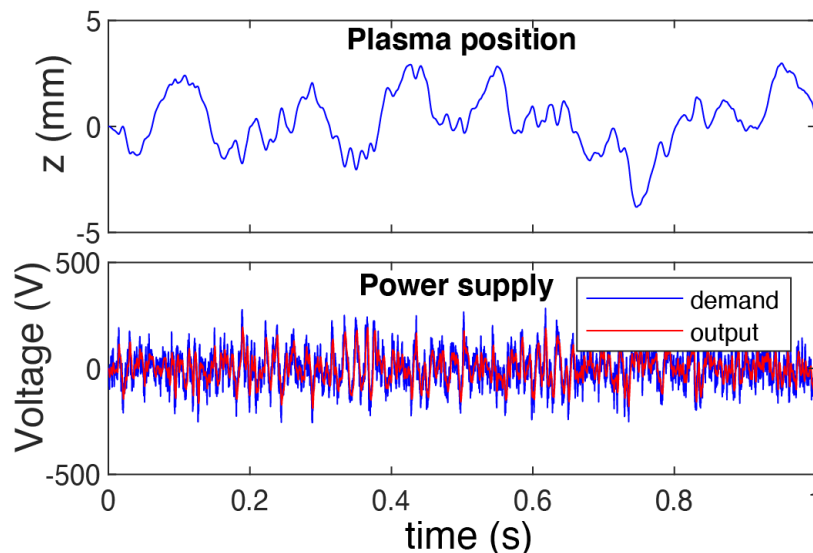
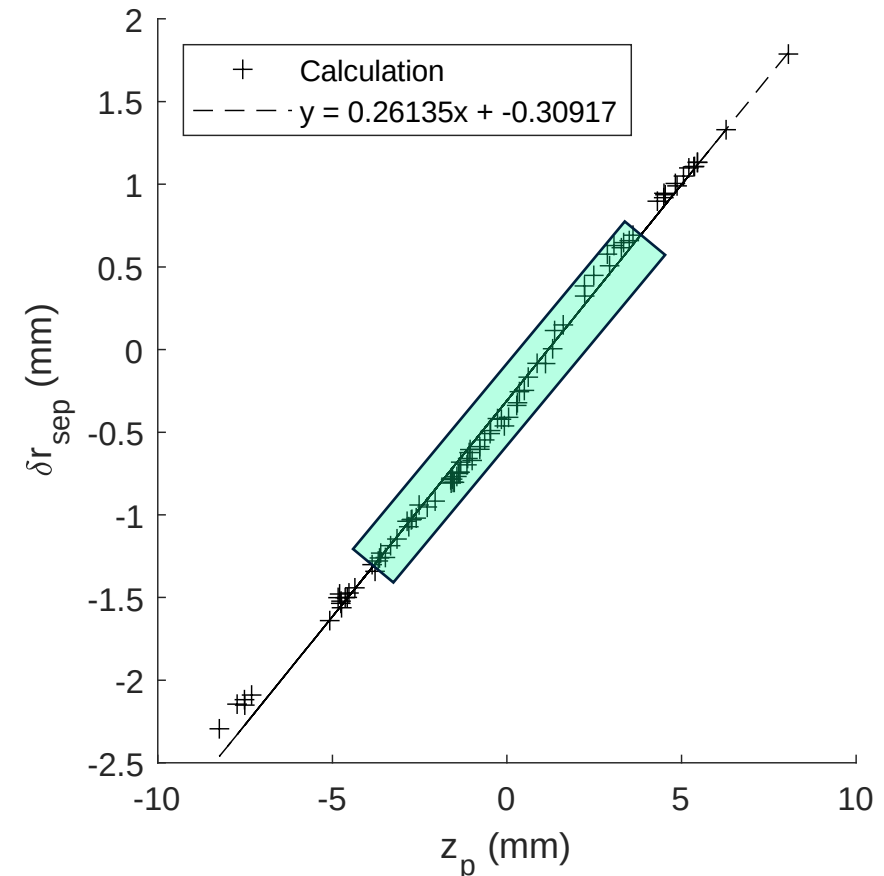
Passive stabiliser rings and in-vessel coils needed for vertical control

- Loss of vertical control can lead to high disruption heat loads
- High accuracy needed to protect inner divertor
 - achieve ± 4 mm with assumed RMS noise of 1 ms^{-1}



Accurate vertical control is needed to protect inner leg by keeping $|\delta r_{sep}| < |\lambda_{SOL}| \sim 1 - 2 \text{ mm}$

- SPP heat loads marginal in single null (SN) $\Rightarrow$ **dynamic** double null (DN) likely needed for commercial plant
- Power distribution between four divertor legs changes dynamically as the plasma oscillates vertically due to control action
- Initial assessment of vertical control system with *mock RMS noise of 1 m/s* shows promising performance
 - In vessel coils + passive stabilising rings



Accuracy of δr_{sep} control in $R = 2.5 \text{ m}$ (SPR-017) device normalised to 1 m/s RMS noise on the vertical velocity measurement.

H&CD mix is pure microwave dimensioned for EC but compatible with EBW

2021 decision:

- Determine HCD mix for STEP ramp-up
- Determine HCD mix for STEP flat-top

2022 decision:

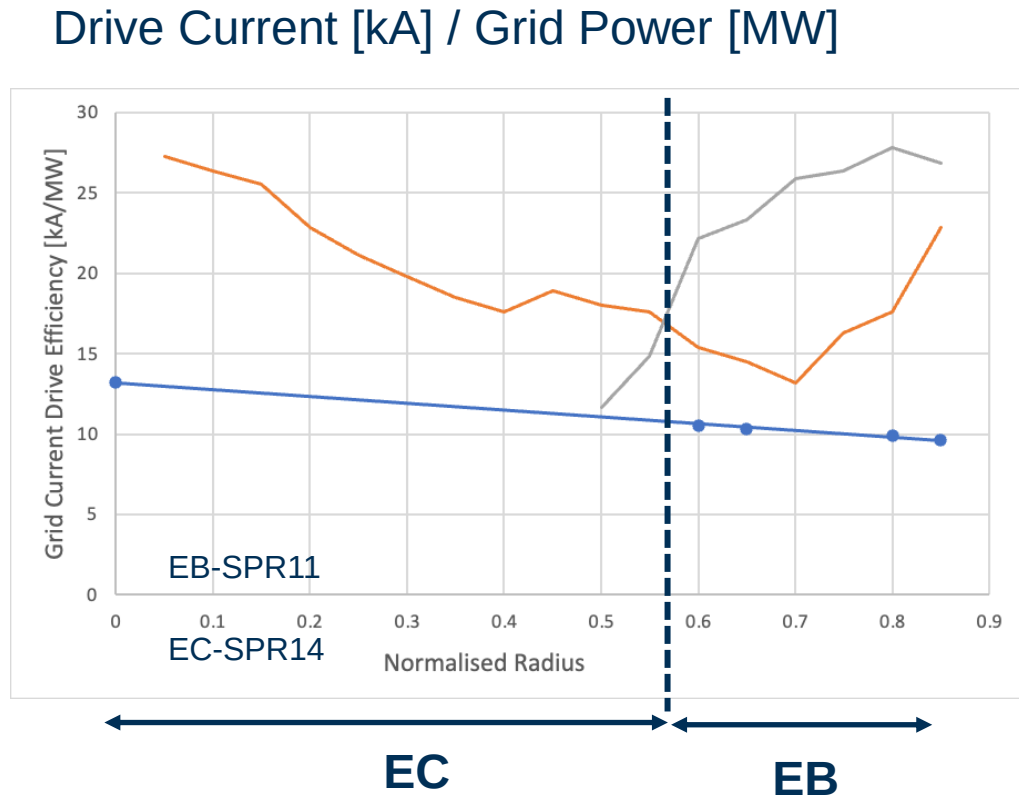
- Determine HCD mix for STEP ramp-down

Options: NB, IC, LH, HW, EB and EC

Documentation:

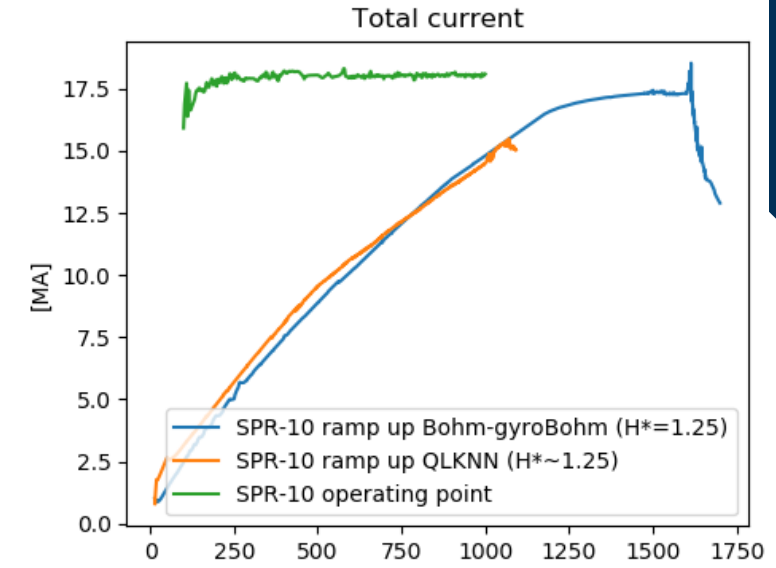
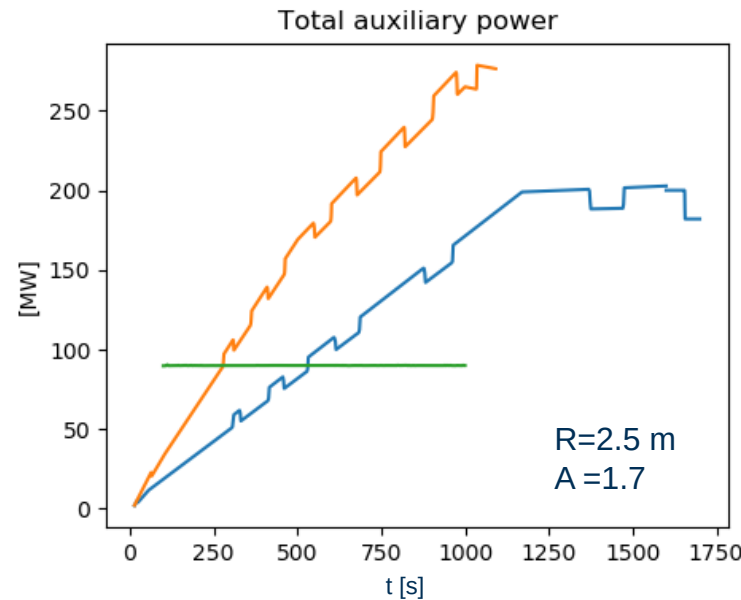
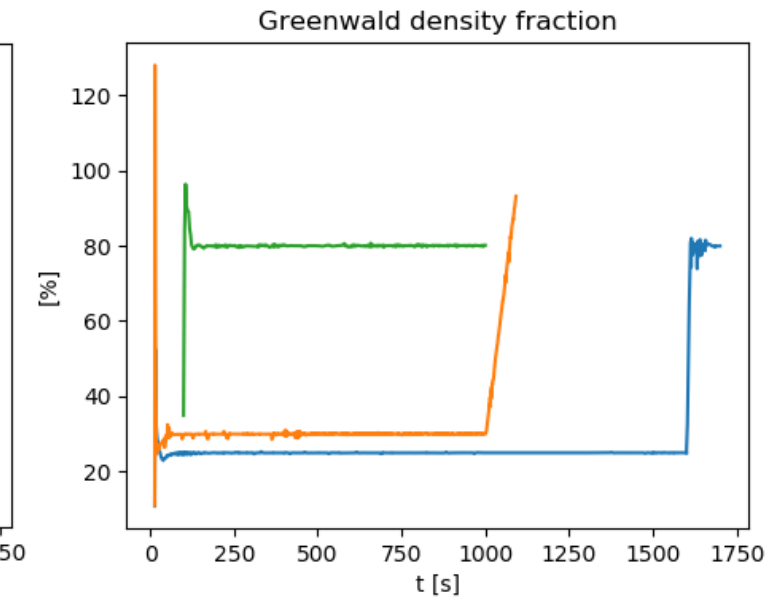
- Define Functional Requirements: [HCD SRD](#)
- Determine evaluation criteria: [HCD Mix Evaluation Criteria](#)
- Evaluation Report: [CML3.5 HCD Mix Evaluation Report](#)

Criteria		NB	EC	EB	HW
1	$\langle \eta_{CD} \rangle$ ($0.5 < \rho < 0.85$) [kA/MW]	~30	~37	~58	TBD
2	Plug-to-Plasma, Target $\eta_{grid} \geq 45\%$	~30%	~44%	~44%	TBD
3	Compatible for Start-up	No	Yes	Partial	No
4	Current Profile Control	Partial	Yes	Partial	TBD
5	VV, Blankets, Magnets Impact	High	Medium	Low	High
6	Building Impact	High	Medium	Medium	TBD
7	System costs	Medium	Low	Low	TBD
8	Impact to STEP Costs	High	Medium	Low	TBD
9	Risks of Technology Level	Medium	Low	High	High
10	RAMI	Low	High	High	Low



Non inductive ramp-up requires high heating & current drive power, & lasts ~ ½ h

- Limited solenoid flux requires non-inductive ramp-up $\Rightarrow$ low density for efficient current drive to reduce P_{CD}
- Access to operating point by “fast” density ramp at full current
- $T_e \gg T_i$ during initial low-density phase.
 - May lead to T_i clamping as observed in AUG & W7-X
- Operating point reached with $P_{CD} > 200 \text{ MW}$.



Non inductive ramp-up requires high heating & current drive power, & lasts $\sim \frac{1}{2}$ h

- Limited solenoid flux requires non-inductive ramp-up $\Rightarrow$ low density for efficient current drive to reduce P_{CD}
- Access to operating point by “fast” density ramp at full current
- $T_e \gg T_i$ during initial low-density phase
 - May lead to T_i clamping as observed in AUG & W7-X
- Operating point reached with $P_{CD} > 200$ MW.

